

# In Pursuit of the Mid-infrared: Modeling pulse propagation in ZBLAN Fiber

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# Acknowledgements

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# Motivation

- We want mid-infrared light ( $3\mu\text{m}$  -  $5\mu\text{m}$ ) primarily because of its applications in absorption spectroscopy.
- There are many gases that absorb light in this range.
- So having a tightly spaced frequency comb in the mid-infrared would be quite nice – allowing us to do things like really sensitive gas detection.

# How to make a pulsed laser at 3 microns

- **Option 1:** Find a 'magical' substance that both lases and efficiently mode-locks around that wavelength.
- **Option 2:** Find a 'magical' fiber that uses non-linear optical effects to turn a regular pulsed laser into a 3 micron pulsed laser.

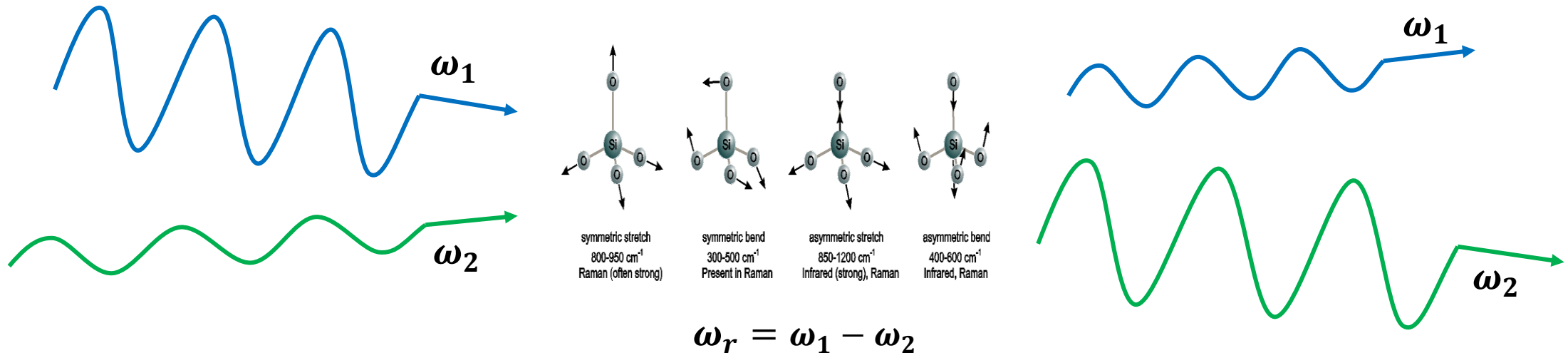
We choose option 2.

# Why does the fiber have to be magical?

- A normal fused silica fiber is also capable of frequency transformation.
- The problem is that for the mid-infrared range, silica is highly absorptive, so doesn't let through much of the light it converts.
- **ZBLAN does not absorb 3 micron light**, providing the needed magic.

# The Frequency Conversion Process

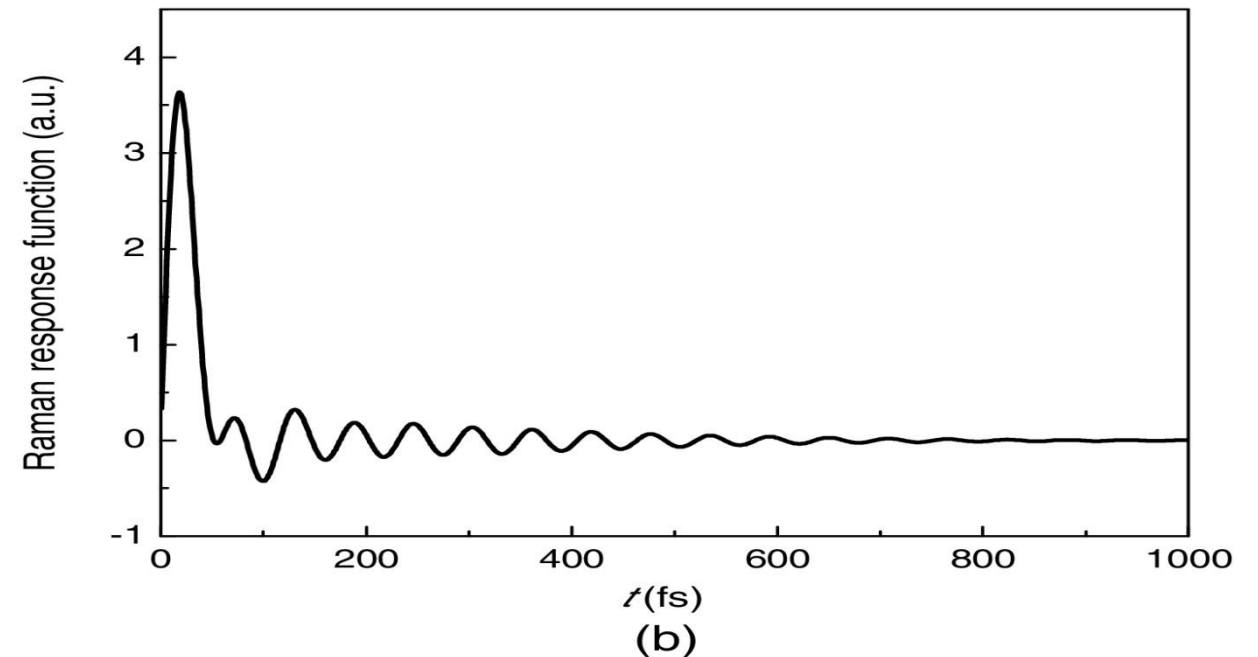
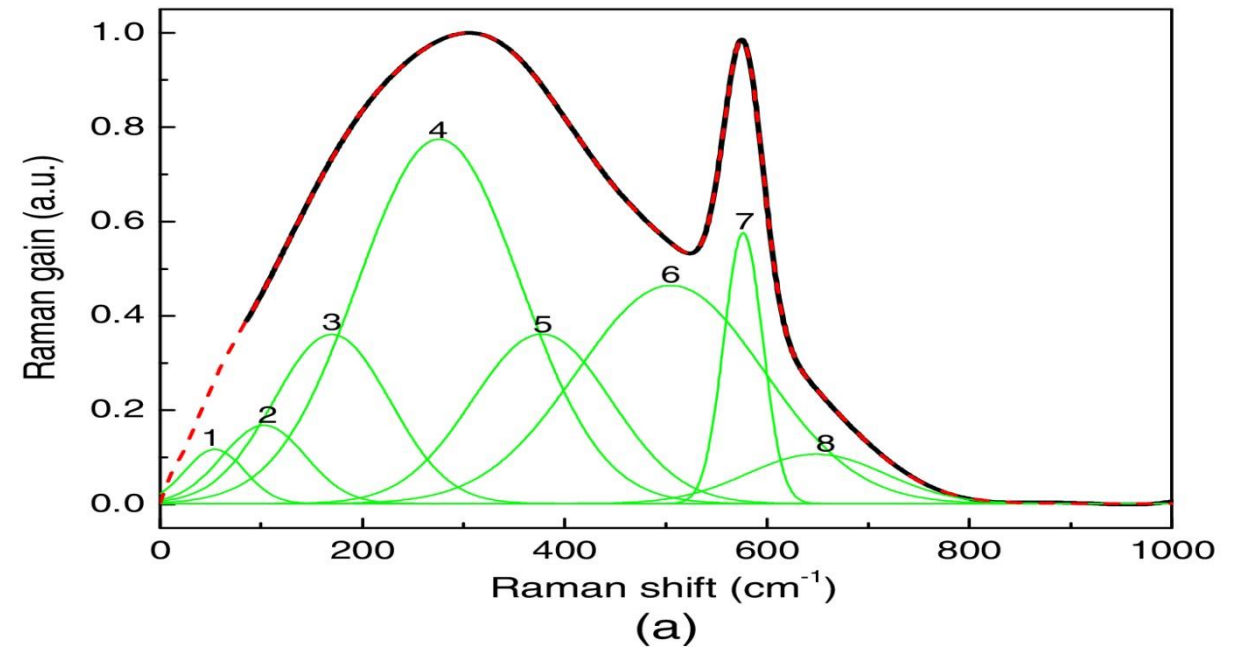
- In a molecular medium, a pair of light waves can resonate with the molecular vibration modes in such a way that energy is transferred from high frequency to low:



- This is called **Stimulated Raman Scattering**

# Raman Scattering in ZBLAN Fiber

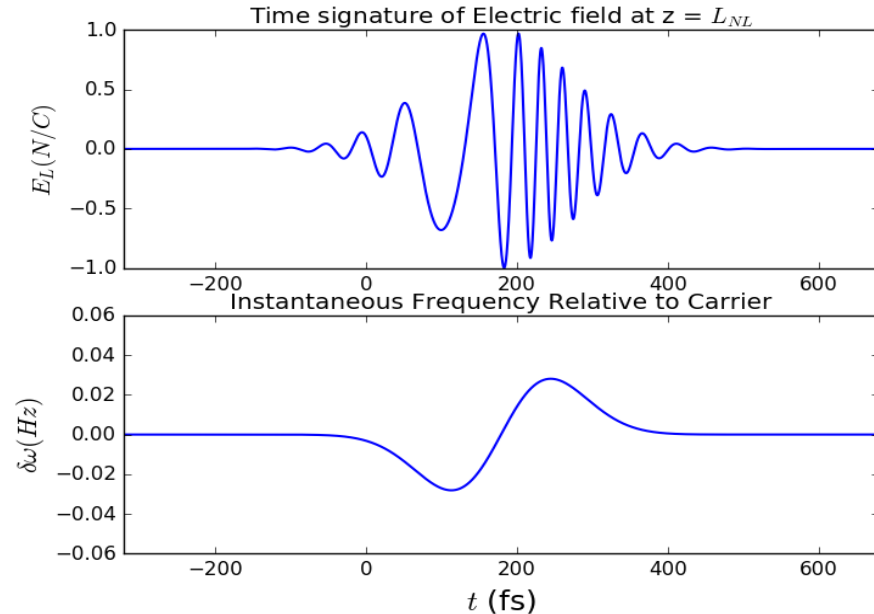
- (a) is a map of the rate at which energy is transferred from one frequency to each frequency below it.
- The black gain curve is a sum of many simple curves that correspond to vibrational modes of ZBLAN molecules.
- The impulse response function (b) corresponding to the Raman gain spectrum tells us everything about how a pulse is affected by Raman scattering.
- **So let's go simulate some pulse evolutions!**



# Not so fast...

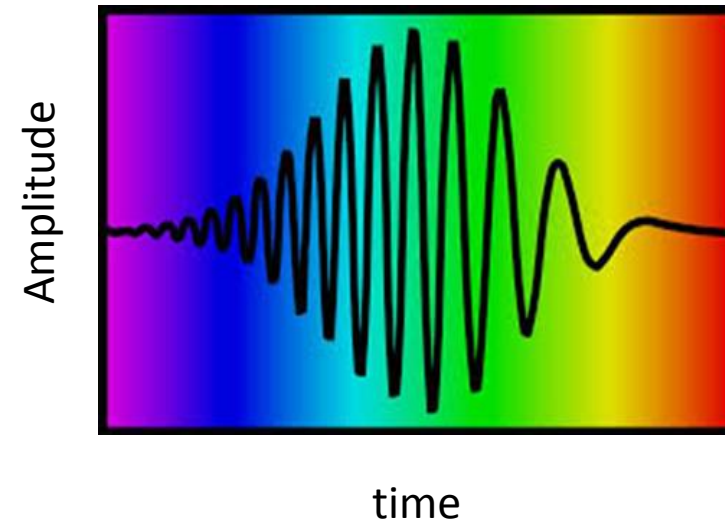
- **There are other effects** that must be incorporated into the simulation to get a realistic picture of pulse propagation through ZBLAN fiber.
- There are two significant ones:

## Self-Phase Modulation



As the pulse travels, the phase of the carrier wave is shifted by an amount proportional to the intensity.

## Dispersion



Different frequency components travel at different group velocities

# How is it possible to consider all these radically different effects at the same time?

- Any optical effect must be describable by Maxwell's wave equation:

$$\nabla^2 E = \mu \frac{\partial^2 (\epsilon_0 E + \vec{P})}{\partial t^2}$$

- The various effects are manifested in the polarization:

$$\vec{P}_{dispersion} = \epsilon_0 \chi_{\omega}^{(1)} \vec{E}$$

$$\vec{P}_{spm} = \epsilon_0 \chi^{(3)} |\vec{E}|^2 \vec{E}$$

$$\vec{P}_{Raman} = \epsilon_0 \chi^{(3)} \vec{E} \int_{-\infty}^{\infty} h_R(t - t') |\vec{E}(t')|^2 dt'$$

# Watering down the Equation

**Approach:** Separate  $\vec{E}$  into its envelope and oscillatory parts and try to get an equation that describes the envelope only.

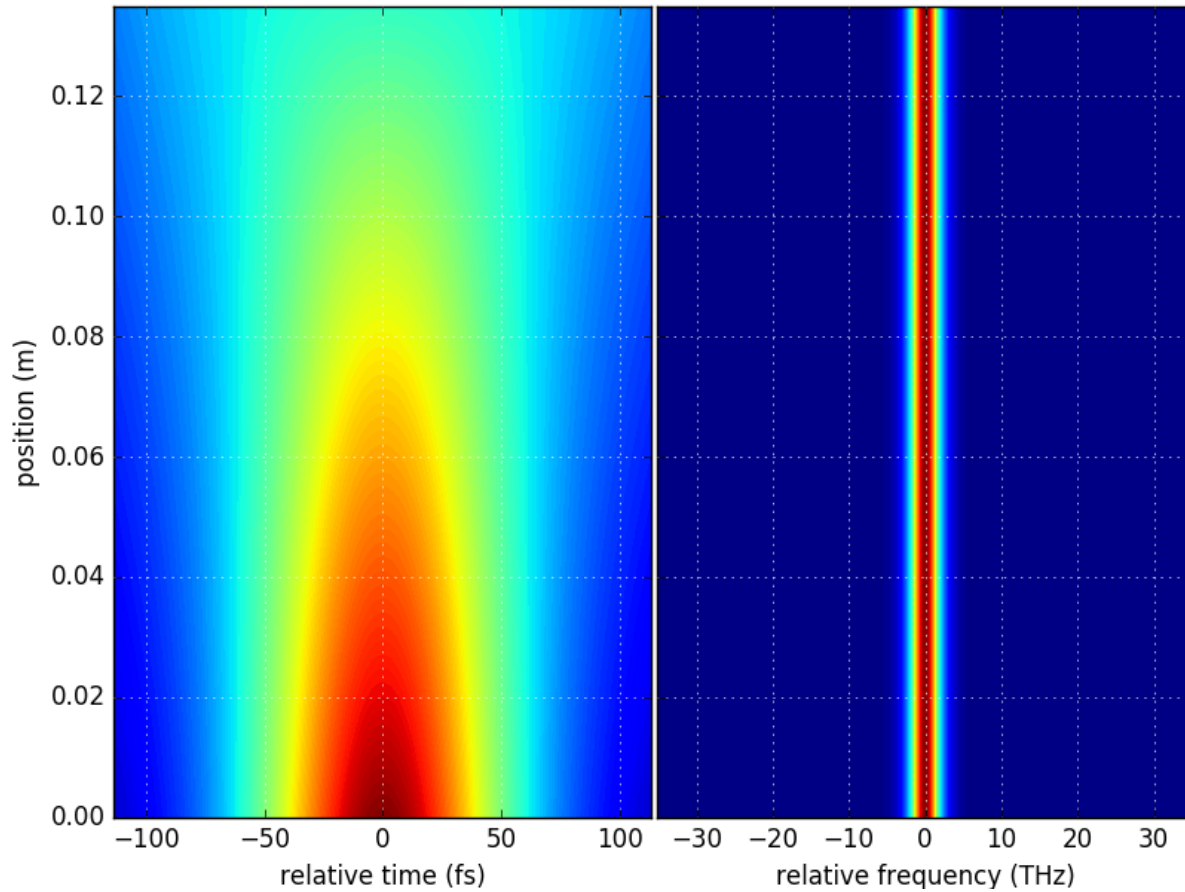
$$E(z, t) = \mathbf{A}(\mathbf{z}, \mathbf{t}) \times e^{i(\omega_0 t - \beta z)}$$

- When this form of  $E$  is plugged into the equation, insignificant second derivatives can be ignored to give a first order ODE of the envelope:

$$\frac{\partial (\mathbf{A}(\mathbf{z}, \omega) e^{-i\beta z})}{\partial z} = i\gamma \frac{\omega}{\omega_0} e^{-i\beta z} FT \quad \mathbf{A}(\mathbf{z}, \mathbf{t}) \int_{-\infty}^{\infty} R(t - t') |\mathbf{A}(\mathbf{z}, \mathbf{t}')|^2 dt'$$

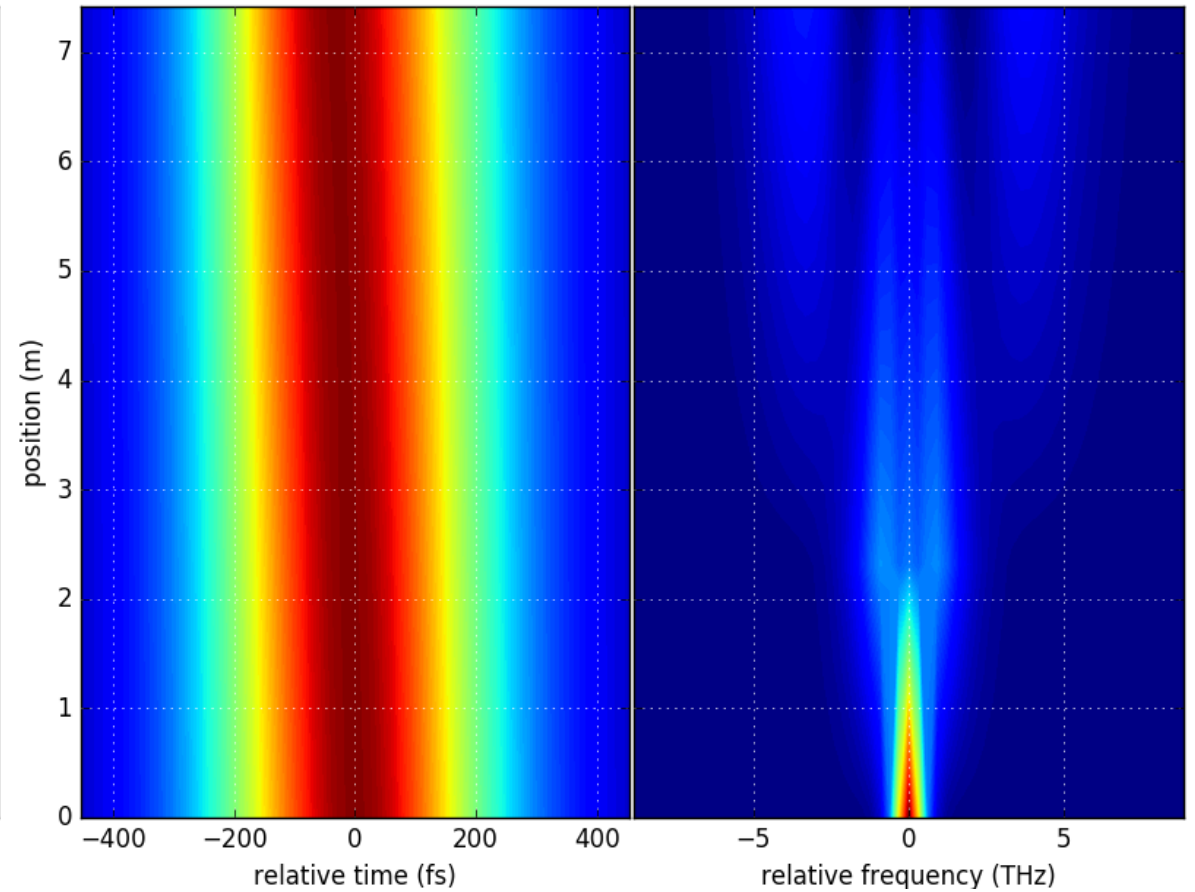
# Solving the envelope equation for each effect individually

$$\vec{P}_{spm}, \vec{P}_{Raman} = 0$$



- As expected, there is no change in spectrum envelope, but broadening in the time envelope.

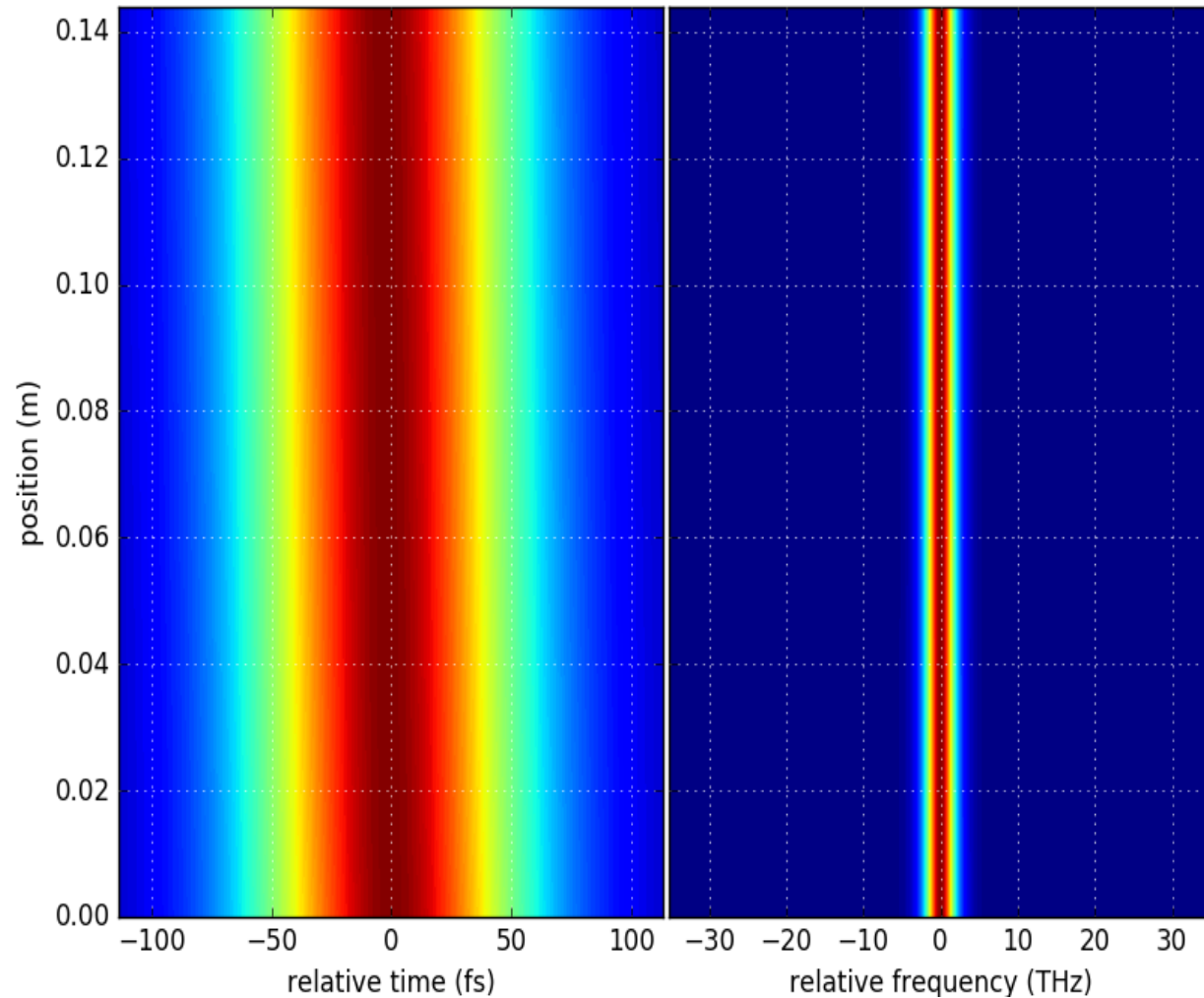
$$\vec{P}_{dispersion}, \vec{P}_{Raman} = 0$$



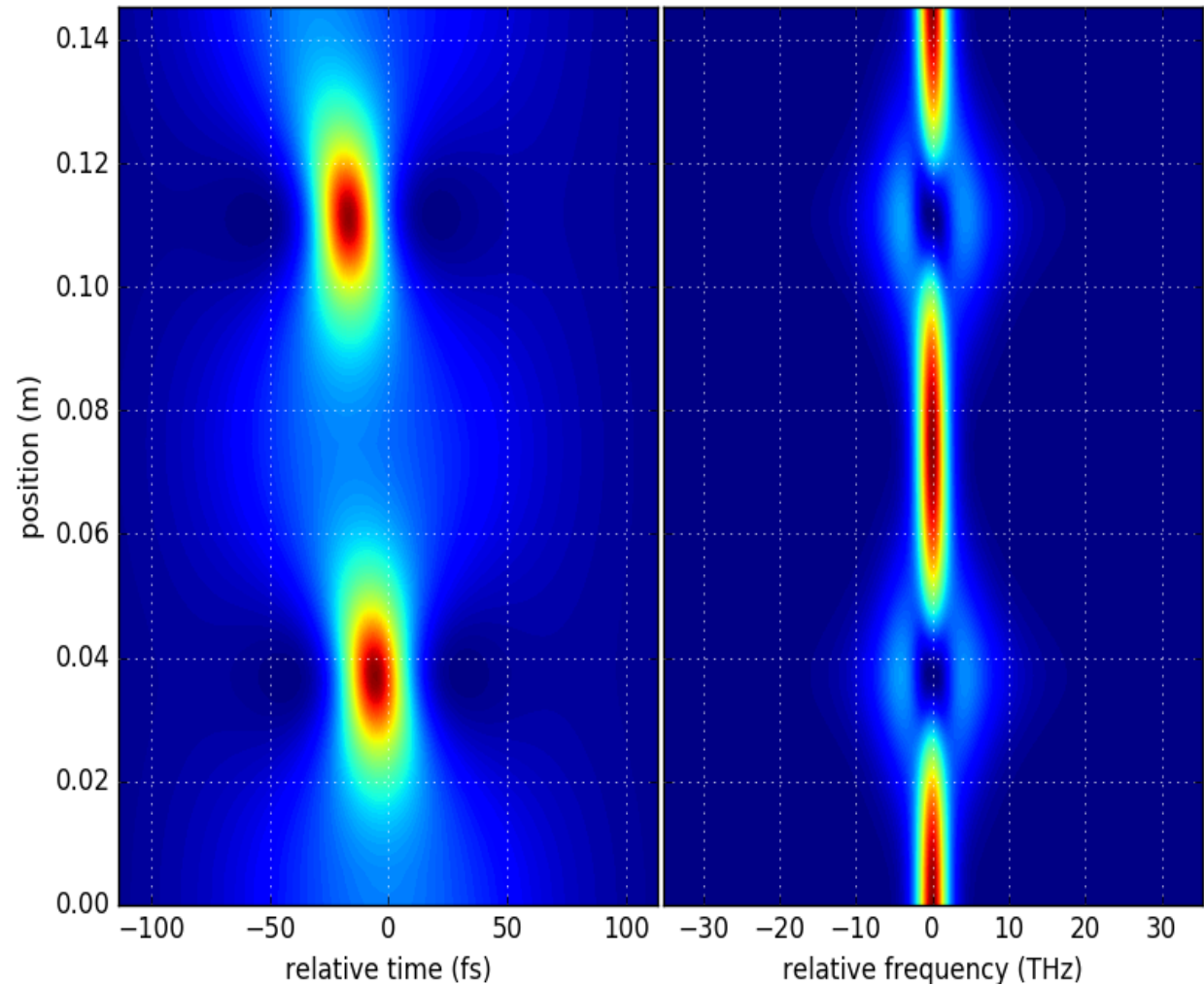
- In the case of self-phase modulation, the spectrum broadens while the time envelope does not change.

# Testing for Solitons (Still no Raman effect)

Soliton order: 1

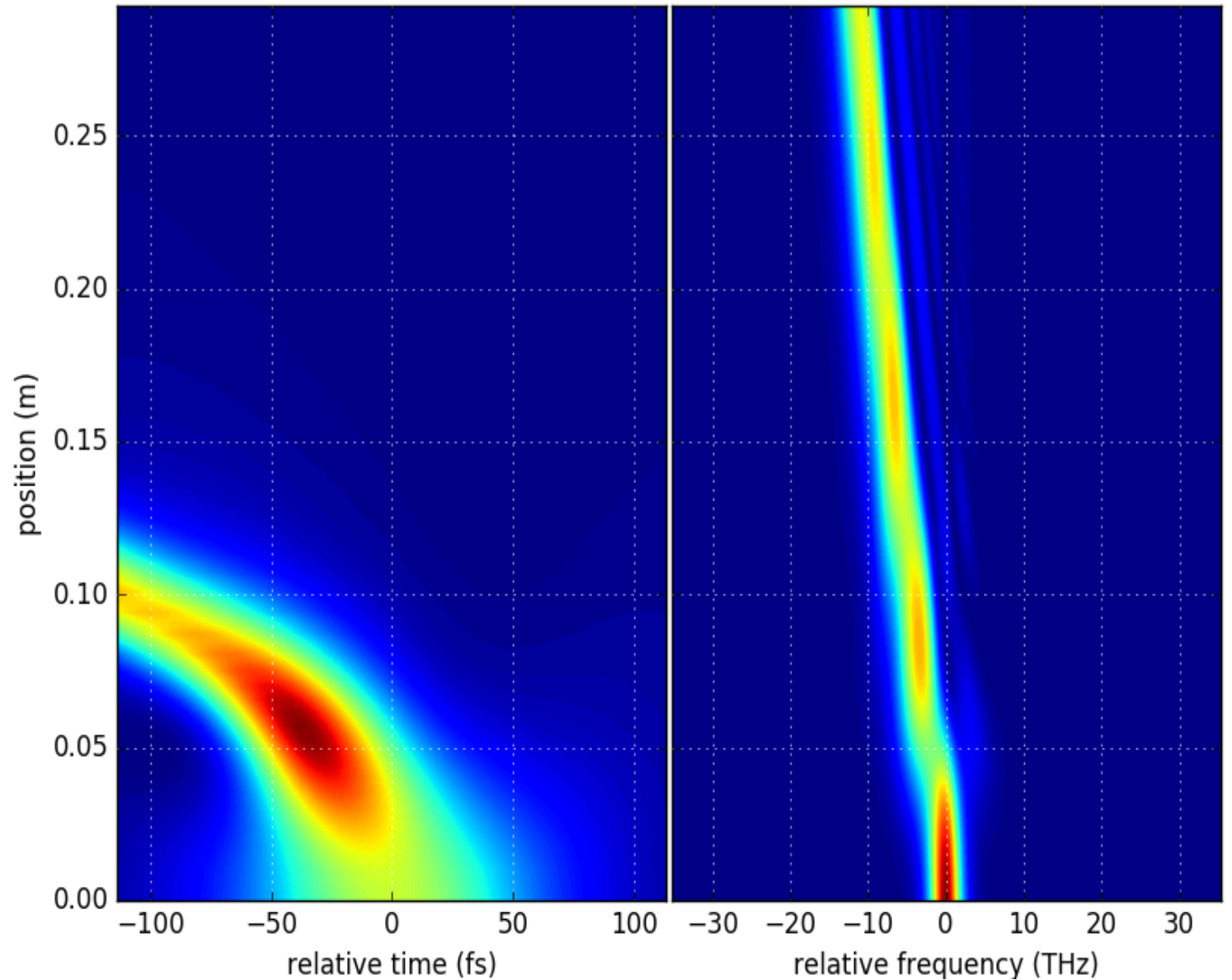


Soliton order: 2



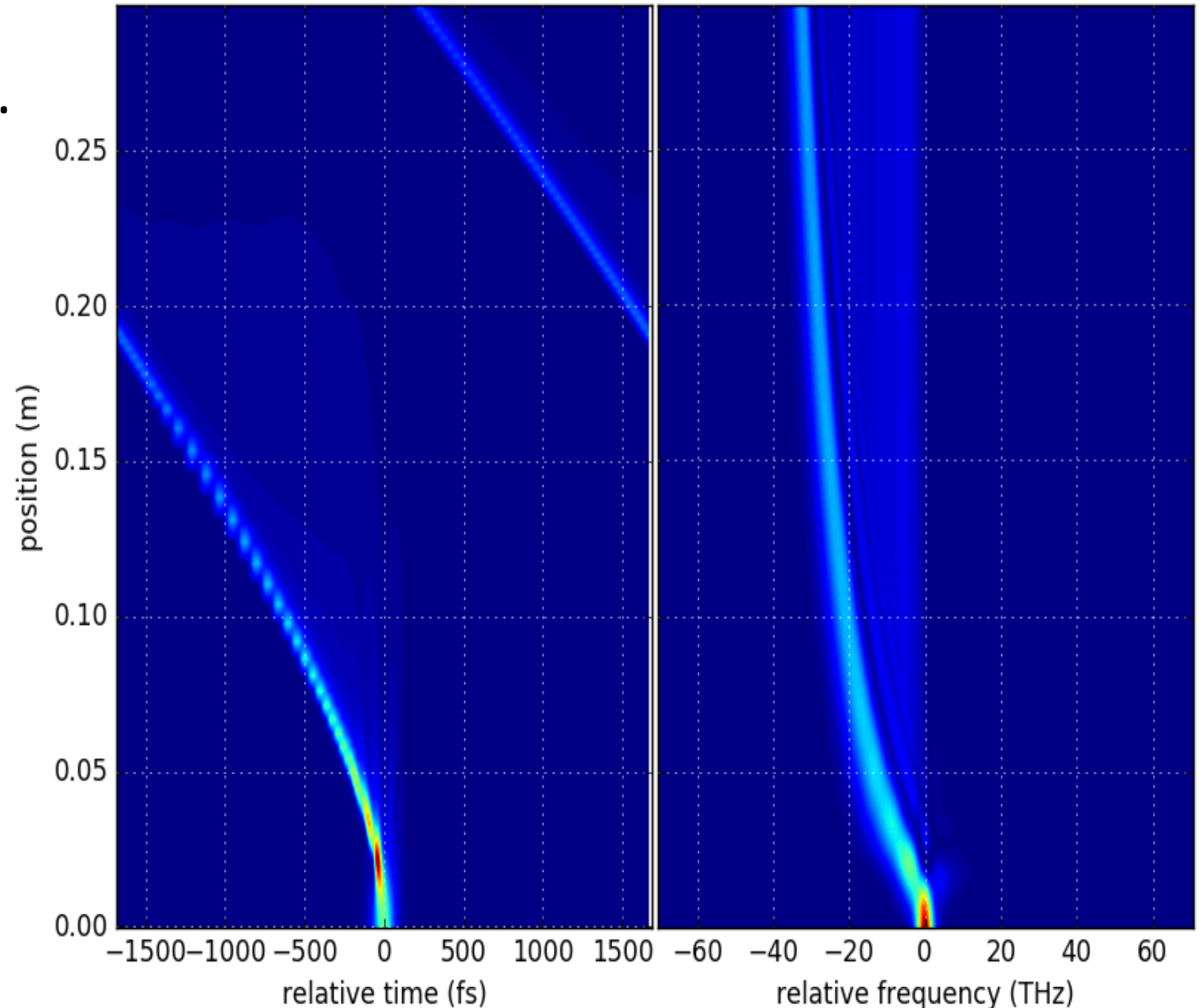
# Working so far, time to include Raman Scattering

- As we were hoping for, a frequency shift is seen when a solitonic pulse undergoes Stimulated Raman Scattering.
- Raman effect is amplified by a factor of 1000 to get results in short propagation distance.



# Pushing the limits...

- Initial pulse is at 2 microns.
- Want to get 3 microns.
- This means that we need a frequency shift of  $c/2\mu\text{m} - c/3\mu\text{m} = 49 \text{ THz}$
- When Raman effect is amplified by another factor of 5, a **negative frequency shift of about 36 THz** is seen.
- Under these conditions, 25% of the initial peak power is retained.



## Conclusions

- The conducted simulations show that ZBLAN fiber is capable of significant frequency shift towards the mid-infrared wavelengths.
- While resonance losses and pulse spreading cannot be neglected, the percentage of transmission appears reasonable.

## Future Works

- Simulate pulse propagation through much longer fiber lengths without enhancing the Raman effect to get more realistic results.
- Look for ideal initial pulse parameters to get as close as possible to the mid-infrared range.

Last Minute!: A negative frequency shift of 50 THz is observed when second order soliton is ejected as the initial pulse.

