

HW #2 solutions

$$\begin{aligned}
 1) \quad S &= -\left(\frac{\partial F}{\partial T}\right)_{N,V} \quad Q = \sum_{E_s} e^{-E_s/kT} \\
 &= k \frac{\partial}{\partial T} (T \ln Q) \\
 &= k \ln Q + \frac{kT}{Q} \frac{\partial Q}{\partial T} \\
 &= k \ln Q + \frac{kT}{Q} \frac{1}{kT^2} \sum E_s e^{-E_s/kT} \\
 &= \frac{k}{Q} \left[Q \ln Q + \frac{1}{kT} \sum E_s e^{-E_s/kT} \right] \\
 &= \frac{k}{Q} \sum_s e^{-E_s/kT} \left[\ln Q + \frac{E_s}{kT} \right] \\
 &= -\frac{k}{Q} \sum_s e^{-E_s/kT} \left[\ln \frac{1}{Q} - \frac{E_s}{kT} \right] \\
 &= -k \sum_s \frac{e^{-E_s/kT}}{Q} \ln \frac{e^{-E_s/kT}}{Q} \\
 &= -k \sum_s f_s \ln \beta_s
 \end{aligned}$$

2

An ideal gas of N particles is confined to a two-dimensional disk of radius R . Each particle is attracted to the center by a force that increases proportional to its distance from the center, so the Hamiltonian is

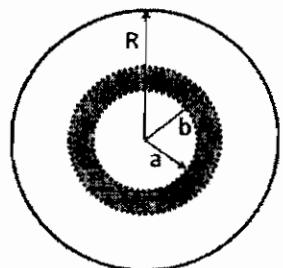
$$H = \sum_{i=1}^N \left(\frac{p_i^2}{2m} + \frac{1}{2} k r_i^2 \right)$$

where k is the effective spring constant of the central force and m is the mass of the particles.

(a) (15 points) Write down the partition function for the full system.

(b) (15 points) Calculate the average energy of the system. *pressure & C_V*

(c) (15 points) Using the single particle partition function, compute the probability that a given particle is in the shaded region of the figure (greater than a distance a from the center but less than b).



a) $Q_N = \frac{Q_1^N}{N!}$

$$= \frac{1}{h^2 N!} \left[\int e^{-\frac{\beta p^2}{2m}} e^{-\frac{\beta kr^2}{2}} dr d\theta \right]^N$$

b) Evaluate $Q_1 = \frac{1}{h^2} \int e^{-\frac{\beta p^2}{2m}} dp \int e^{-\frac{\beta kr^2}{2}} dr$

$$= \frac{1}{h^2} \frac{2m\pi}{\beta} \int_0^{2\pi} \int_0^R e^{-\beta k R^2/2} R dR \quad x = \frac{\beta k R^2}{2}$$

$$dx = \beta k R dR$$

$$= \frac{4m\pi^2}{\beta^2 h^2} \int_0^{\beta k R^2/2} e^{-x} \frac{dx}{\beta k}$$

$$= \frac{4m\pi^2}{\beta^2 k h^2} \left[-e^{-\beta k R^2/2} + 1 \right]$$

$$\langle E \rangle = -\frac{\partial \ln Q_1}{\partial \beta} = -\frac{\partial}{\partial \beta} \left[\ln Q_1 - \ln N! \right] \xrightarrow{\text{constant}}$$

$$= +\frac{2}{\partial \beta} \left[N \ln \frac{\beta^2 k h^2}{4m\pi^2} - N \ln \left(1 - e^{-\beta k R^2/2} \right) \right] \xrightarrow{\frac{2}{\beta} = 2kT(N)}$$

$$= \frac{2N}{\beta} - N \frac{\ln^2 \frac{k R^2}{2} e^{-\beta k R^2/2}}{1 - e^{-\beta k R^2/2}} \xrightarrow[k \rightarrow 0]{R \rightarrow \infty} \frac{2N}{\beta} - \frac{\frac{k R^2}{2} N}{1 - (1 - \frac{k R^2}{2})} = \frac{N = k T N}{\beta} \xrightarrow[\text{equipartition}, c.]{}$$

$$\text{Pressure } P = -\frac{\partial F}{\partial V}$$

in this case we are in two dimensions so "Volume" is really the area = πR^2

$$P = -\frac{\partial F}{\partial A} = \frac{\partial R}{\partial A} \frac{\partial F}{\partial R}$$

$$= -\frac{1}{2\pi R} \frac{\partial F}{\partial R}$$

$$A = \pi R^2$$

$$\frac{\partial A}{\partial R} = 2\pi R$$

$$F = -kT \ln Q_N = -kT \ln \frac{Q_N}{N!} = -NkT \ln Q + NkT \ln N - NkT N$$

$$= -NkT \ln \frac{Q}{N} - NkT$$

$$= -NkT \ln \left[\frac{4\pi R^2}{N\beta^2 k h^2} \left(1 - e^{-\beta k R^2/2} \right) \right] - NkT$$

$$= -NkT \ln \left[\frac{\lambda}{\beta k N \lambda^2} \left(1 - e^{-\beta k R^2/2} \right) \right] - NkT \quad \lambda = \frac{\hbar}{\sqrt{2m\pi k_B T}}$$

$$\frac{\partial E}{\partial R} = \frac{\partial kR}{\partial R} \frac{e^{-\beta k R^2/2}}{1 - e^{-\beta k R^2/2}} (-NkT)$$

$$P = -\frac{1}{2\pi R} \left(-NkT \right) \frac{\lambda k R}{e^{\beta k R^2/2} - 1} = \boxed{\frac{Nk}{2\pi \left(e^{\beta k R^2/2} - 1 \right)}}$$

$$\lim_{k \rightarrow 0} k \rightarrow 0$$

$$P \rightarrow \frac{Nk}{2\pi \left(\lambda + \beta k R^2/2 - 1 \right)}$$

$$\therefore P(\pi R^2) = Nk_B T \quad \checkmark$$

$$C_V = \frac{\partial E}{\partial T} = \frac{\partial E}{\partial T} \frac{\partial E}{\partial \beta}$$

$$\frac{\partial \beta}{\partial T} = -\frac{1}{kT^2}$$

$$= -\frac{1}{kT^2} \left[-\frac{2N}{\beta^2} - \frac{NkR^2}{2} \frac{2}{\partial \beta} \left(\frac{1}{e^{\beta k R^2/2} - 1} \right) \right]$$

$$= -\frac{\beta}{T} \left[-\frac{2N}{\beta^2} - \frac{NkR^2}{2} \frac{-\left(kR^2/2 \right) e^{\beta k R^2/2}}{\left(e^{\beta k R^2/2} - 1 \right)^2} \right]$$

$$= \frac{\beta}{T} \left[\frac{2N}{\beta^2} - N \left(\frac{kR^2}{2} \right)^2 \frac{e^{\beta k R^2/2}}{\left(e^{\beta k R^2/2} - 1 \right)^2} \right] \quad \leftarrow$$

$$C) Q_{ab} = \frac{1}{h^2} \underbrace{\int_a^b e^{-\beta k R^{3/2}} R dR}_{\text{momentum integral}}$$

$$\text{Probability} = \frac{Q_{ab}}{Q_i} = \frac{\int_a^b e^{-\beta k R^{3/2}} R dR}{\int_0^R e^{-\beta k R^{3/2}} R dR}$$

$$= \boxed{\frac{e^{-\beta k a^{3/2}} - e^{-\beta k b^{3/2}}}{1 - e^{-\beta k R^{3/2}}}}$$

P 3.15

$$Q_N = \frac{1}{N!} Q^N \quad Q = \frac{1}{h^3} \int e^{-\beta p c} d^3 p d^3 q$$

$$= \frac{4\pi V}{h^3} \int e^{-\beta p c} p^2 dp$$

$$x = \beta p c$$

$$dx = \beta c dp$$

$$u = x^2 \quad v = -e^{-x}$$

$$du = 2x dx \quad dv = e^{-x} dx$$

$$= \frac{4\pi V}{(\beta h c)^3} \int e^{-x} x^2 dx$$

$$= x^2 e^{-x} \Big|_0^\infty + 2 \int x e^{-x} dx$$

$$- x e^{-x} \Big|_0^\infty$$

$$= \frac{8\pi V}{(\beta h c)^3}$$

$$Q_N = \frac{1}{N!} \left(\frac{8\pi V}{(\beta h c)^3} \right)^N$$

$$= -k T \ln \left(\frac{8\pi V}{N} \left(\frac{k T}{h c} \right)^3 + 1 \right)$$

$$P = - \left(\frac{\partial F}{\partial V} \right) = \frac{N k T}{V}$$

$$E = - \frac{\partial}{\partial \beta} \ln Q = \frac{\partial}{\partial \beta} [3N \ln \beta + \text{constant}]$$

$$E = 3N k T$$

$$PV = N k T = \frac{E}{3}$$

$$C_V = \frac{\partial E}{\partial T} = 3N k$$

$$C_P = \frac{\partial (E + PV)}{\partial T} = \frac{2}{T} (3N k T + N k T) \\ = 4N k$$

$$\gamma = \frac{C_P}{C_V} = \frac{4N k}{3N k} = \frac{4}{3}$$

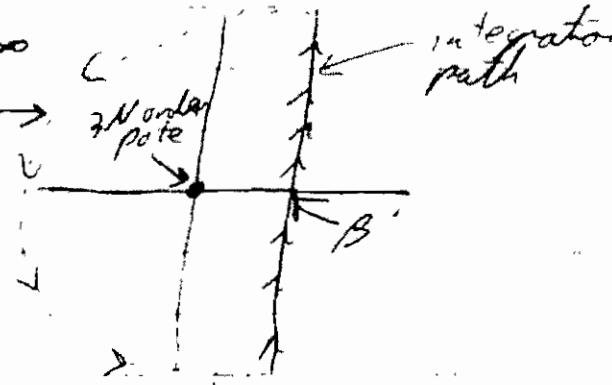
Pathria 3.5
(cont.)

$$g(E) = \frac{1}{2\pi i} \int_{\beta=-\infty}^{\beta=\infty} e^{\beta E} Q_n(\beta) d\beta$$

$$= \frac{i}{2\pi i} \frac{1}{N!} \left(\frac{8\pi V}{(hc)^3} \right)^N \int \frac{e^{\beta E}}{\beta^{3N}} d\beta$$

Close contour with $R=\infty$

loop in Right half plane \rightarrow left



Loop does not contribute
d/dt due to factor β^{-3N}

(note: exponential factor is

bounded $< e^{\beta E}$ due to closure on left)

$$\oint e^{\beta E} \beta^{-3N} d\beta = 2\pi i \underbrace{\text{Res}(\beta=0)}_{}$$

$$\lim_{A \rightarrow 0} \frac{1}{(3N-1)!} E^{(3N-1)} e^{\beta E}$$

$$g(E) = \frac{E^{3N-1}}{N! (3N-1)!} \left(\frac{8\pi V}{(hc)^3} \right)^N$$

4) a) N dipoles N_+ with energy ε
 N_- with energy $-\varepsilon$ $N_+ + N_- = N$

$$E = N_+ \varepsilon - N_- \varepsilon = (N_+ - N_-) \varepsilon = (2N_+ - N) \varepsilon$$

$$N_+ = \frac{E/\varepsilon + N}{2} \quad N_- = N - N_+ = \frac{N - E/\varepsilon}{2}$$

of states

$$\Omega = \binom{N}{N_+} = \frac{N!}{N_+! N_-!}$$

$$\frac{S}{k_B} = \ln \Omega$$

$$\approx \underbrace{N \ln N}_{(N_+ + N_-) \ln N} - N_+ \ln N_+ - N_- \ln N_- + N_+ \ln N_+ + N_- \ln N_-$$

$$= -N_+ \ln \frac{N_+}{N} - N_- \ln \frac{N_-}{N}$$

$$\frac{S}{Nk_B} = -\frac{N_+}{N} \ln \frac{N_+}{N} - \frac{N_-}{N} \ln \frac{N_-}{N}$$

$$\rightarrow \frac{N_+}{N} = \frac{N\varepsilon + E}{2N\varepsilon} \quad \frac{N_-}{N} = \frac{N\varepsilon - E}{2N\varepsilon}$$

$$\frac{S}{Nk_B} = -\frac{N\varepsilon + E}{2N\varepsilon} \ln \frac{N\varepsilon + E}{2N\varepsilon} - \frac{N\varepsilon - E}{2N\varepsilon} \ln \frac{N\varepsilon - E}{2N\varepsilon}$$

b) Canonical version

$$Q = \sum_E g(E) e^{-\beta E}$$

or $\sum_{N_+}^N g(N_+) e^{-\beta(N_+ \epsilon - N_- \epsilon)}$

$$= \sum_{N_+} g(N_+) e^{-\beta(2N_+ - N) \epsilon}$$

$$= e^{N\beta\epsilon} \sum_{N_+=0}^N \binom{N}{N_+} e^{-\beta 2N_+ \epsilon}$$

$$= e^{N\beta\epsilon} (1 + e^{-2\beta\epsilon})^N$$

$$= (e^{\beta\epsilon} + e^{-\beta\epsilon})^N$$

$$= Q_1^N$$

E goes from $-N\epsilon$ to $N\epsilon$
in steps of 2ϵ

$$g(N_+) = \binom{N}{N_+} = \frac{N!}{N_+!(N-N_+)!}$$

← Binomial Series!

Answers & Solutions

1) a)

$$\begin{aligned}
 Q_1 &= \sum_{m=-j}^j e^{-\beta \mu g H m} = \sum_{-j}^j e^{-mx} \quad x = \beta \mu g H \\
 &= e^{xj} \sum_{n=0}^{2j} e^{-nx} \\
 &= e^{xj} \left[\sum_{n=0}^{\infty} e^{-nx} - \sum_{n=2j+1}^{\infty} e^{-nx} \right] \\
 &= e^{xj} \left[\frac{1}{1-e^{-x}} - e^{-x(2j+1)} \frac{1}{1-e^{-x}} \right] \\
 &= \frac{e^{xj}}{1-e^{-x}} \left[1 - e^{-x(2j+1)} \right] = \frac{e^{xj} - e^{-x(j+1)}}{e^{-x/2} (e^{x/2} - e^{-x/2})} = \boxed{\frac{\sinh(xj + \frac{x}{2})}{\sinh(\frac{x}{2})}}
 \end{aligned}$$

b)

$$\begin{aligned}
 M_2 &= \frac{\partial \ln Q}{\partial (\beta H)} = \frac{\partial x}{\partial \beta H} \frac{\partial}{\partial x} \left[\ln \left\{ \sinh \left(x \left(j + \frac{1}{2} \right) \right) \right\} - \ln \left(\frac{i \tanh \frac{x}{2}}{2} \right) \right] \\
 &= \mu g \left[\left(j + \frac{1}{2} \right) \coth \left\{ x \left(j + \frac{1}{2} \right) \right\} - \frac{1}{2} \coth \left(\frac{x}{2} \right) \right]
 \end{aligned}$$

c)

$$\begin{aligned}
 M_2 &= \mu g \left[\coth(x) - \frac{1}{2} \coth \left(\frac{x}{2} \right) \right] \\
 &= \mu g \left[\frac{2 \cosh x \sinh \frac{x}{2} - \cosh \frac{x}{2} \sinh x}{2 \sinh x \sinh \frac{x}{2}} \right] \\
 &= \mu g \left[\frac{2 \left(\sinh \frac{x}{2} \cosh^2 \frac{x}{2} + \sinh^3 \frac{x}{2} \right) - 2 \cosh^2 \frac{x}{2} \sinh^2 \frac{x}{2}}{2 \sinh^2 \frac{x}{2} \cosh \frac{x}{2}} \right] \\
 &= \mu g \frac{\sinh \frac{x}{2}}{\cosh \frac{x}{2}} = \mu g \tanh(\beta \mu g H) \\
 &\text{equivalent to lecture of } g=2
 \end{aligned}$$

C (cont.)

$$J \rightarrow \infty \quad (g \rightarrow 0)$$

$$\begin{aligned} M_2 &= \mu g \left[J \coth \left\{ x \left(J + \frac{1}{2} \right) \right\} + \frac{1}{2} \cancel{\mu g} \overbrace{\coth \left\{ x \left(J + \frac{1}{2} \right) \right\}}^0 - \frac{\mu g}{2} \underbrace{\coth \left(\frac{x}{2} \right)}_{\substack{\coth x \rightarrow \frac{1}{x} \\ \text{as } x \rightarrow 0}} \right] \\ &= \mu_0 \coth \left\{ \beta \mu g H + \frac{1}{2} \cancel{\beta \mu g H} \right\} - \frac{\mu_0}{x} \\ &= \mu_0 \coth \left(\beta H \mu_0 \right) - \frac{1}{\beta H} \quad \leftarrow \text{classical result} \end{aligned}$$

d) $\lim_{x \rightarrow 0} M_2 = \mu g \left[\left(J + \frac{1}{2} \right) \cancel{\frac{1}{x \left(J + \frac{1}{2} \right)}} \left(1 + \frac{1}{3} x^2 \left(J + \frac{1}{2} \right)^2 \right) - \frac{1}{2} \cancel{\frac{2}{x}} \left(1 + \frac{1}{3} \left(\frac{x}{2} \right)^2 \right) \right]$

$$\begin{aligned} &= \mu g \left[\frac{1}{x} \left(J + \frac{1}{3} x^2 \left(J + \frac{1}{2} \right)^2 \right) - \frac{1}{x} \left(J + \left(\frac{x}{2} \right)^2 \right) \right] \\ &= \cancel{\mu g} \frac{x}{3} \left(J^2 + J \right) \end{aligned}$$

Curie constant is : $\frac{(\mu g)^2 J(J+1)}{k_B T}$