

Chapter 6, Time-independent perturbation Theory

1. Nondegenerate perturbation theory

To solve $H\Psi = E\Psi$

$$\text{but } H = H^0 + H'$$

suppose we know how to solve H^0 , and H' is small,
we want to obtain approximate solution of $H\Psi = E\Psi$

$$\text{Let } H = H^0 + \lambda H'$$

Assume we know

$$H^0 \varphi_n^0 = E_n^0 \varphi_n^0$$

$$\text{Write } H \varphi_n = E_n \varphi_n$$

$$\text{Expand } \varphi_n = \varphi_n^0 + \lambda \varphi_n^1 + \lambda^2 \varphi_n^2 + \dots$$

$$E_n = E_n^0 + \lambda E_n^1 + \lambda^2 E_n^2 + \dots$$

plug in

$$\begin{aligned} & (H^0 + \lambda H') (\varphi_n^0 + \lambda \varphi_n^1 + \lambda^2 \varphi_n^2 + \dots) \\ &= (E_n^0 + \lambda E_n^1 + \lambda^2 E_n^2 + \dots) (\varphi_n^0 + \lambda \varphi_n^1 + \lambda^2 \varphi_n^2 + \dots) \end{aligned}$$

Collecting terms of different orders in λ

$$\lambda^0 \quad H^0 \varphi_n^0 = E_n^0 \varphi_n^0 \tag{1}$$

$$\lambda^1 \quad H^0 \varphi_n^1 + H' \varphi_n^0 = E_n^0 \varphi_n^1 + E_n^1 \varphi_n^0 \tag{2}$$

$$\lambda^2 \quad H^0 \varphi_n^2 + H' \varphi_n^1 = E_n^0 \varphi_n^2 + E_n^1 \varphi_n^1 + E_n^2 \varphi_n^0 \tag{3}$$

Note that $\{\varphi_n^0\}$ orthonormal
and that it is not degenerate.

From (2), take $\langle \varphi_n^0 |$ on both sides

$$\langle \varphi_n^0 | H^0 | \varphi_n^0 \rangle + \langle \varphi_n^0 | H' | \varphi_n^0 \rangle = E_n^0 \langle \varphi_n^0 | \varphi_n^0 \rangle + E_n^1$$

\downarrow ←
 $E_n^0 \langle \varphi_n^0 | \varphi_n^0 \rangle \leftarrow$ cancell

Thus
$$\boxed{E_n^1 = \langle \varphi_n^0 | H' | \varphi_n^0 \rangle} \quad (4)$$

\therefore 1st order perturbation to the energy E_n^0

Note: To calculate 1st order correction to energy, we need only φ_n^0 .

From (2), we solve

$$(H^0 - E_n^0) \varphi_n^0 = -(H' - E_n^1) \varphi_n^0$$

Note that we can expand φ_n^0 in the basis of $\{\varphi_m^0\}$

$$\varphi_n^0 = \sum_{m \neq n} c_m \varphi_m^0$$

$m=n$ is excluded

$$(H^0 - E_n^0) \sum_{m \neq n} c_m \varphi_m^0 = -(H' - E_n^1) \varphi_n^0$$

Take $\langle \varphi_m^0 |$ from the left $(m \neq n)$

$$(E_m^0 - E_n^0) c_m = - \langle \varphi_m^0 | H' | \varphi_n^0 \rangle$$

$$c_m = \frac{\langle \varphi_m^0 | H' | \varphi_n^0 \rangle}{E_n^0 - E_m^0}$$

or
$$\boxed{\varphi_n^1 = \sum_{m \neq n} \frac{\langle \varphi_m^0 | \varphi_n^0 \rangle \langle \varphi_m^0 | H' | \varphi_n^0 \rangle}{E_n^0 - E_m^0}} \quad (5)$$

1st order correction to the wavefunction

2nd-order correction of energy-

From Eq. (3)

$$\langle \varphi_n^0 | H^0 | \varphi_n^0 \rangle + \langle \varphi_n^0 | H' | \varphi_n^0 \rangle = E_n^0 \cancel{\langle \varphi_n^0 | \varphi_n^2 \rangle} + E_n^1 \langle \varphi_n^0 | \varphi_n^1 \rangle + E_n^2$$

Since $\langle \varphi_n^1 \rangle = \sum_{m \neq n} \epsilon_m |\varphi_m^0 \rangle$

Note that $\langle \varphi_n^0 | \varphi_n^1 \rangle = 0$ from the eq. above

Thus

$$E_n^2 = \langle \varphi_n^0 | H' | \varphi_n^1 \rangle = \langle \varphi_n^0 | H' | \sum_{m \neq n} \frac{|\varphi_m^0 \rangle \langle \varphi_m^0 | H' | \varphi_n^0 \rangle}{E_n^0 - E_m^0} \rangle$$

$$E_n^2 = \sum_{m \neq n} \frac{|\langle \varphi_m^0 | H' | \varphi_n^0 \rangle|^2}{E_n^0 - E_m^0}$$

↑
2nd order correction to energy

General comments:

- (1) No need to go to E_n^2 if E_n^1 is not zero. In other words, you need to calculate only the first nonzero term.
- (2) The calculation of E_n^2 is not trivial since you need to sum over all the states (except $m=n$) in the Hilbert space.
- (3) Many matrix elements need to be evaluated.

2. Degenerate perturbation Theory

For simplicity, consider the case that H^0 has two-fold degeneracy.

$$H^0 \varphi_1^0 = E^0 \varphi_1^0$$

$$H^0 \varphi_2^0 = E^0 \varphi_2^0$$

We can assume $\langle \varphi_1^0 | \varphi_2^0 \rangle = 0$

We look for solutions

$$H \Psi = E \Psi$$

Where $H = H^0 + H'$

We will look for

$$\Psi = a \varphi_1^0 + b \varphi_2^0 + \dots$$

but neglecting the other terms that are not degenerate with E^0 .

From $(H^0 + H')(a \varphi_1^0 + b \varphi_2^0) = E(a \varphi_1^0 + b \varphi_2^0)$

$$\begin{aligned} a \cdot H^0 \varphi_1^0 + b \cdot H^0 \varphi_2^0 + a \cdot H' \varphi_1^0 + b \cdot H' \varphi_2^0 \\ = a E^0 \varphi_1^0 + b E^0 \varphi_2^0 \end{aligned}$$



$$\begin{aligned} \langle \varphi_1^0 | & a E^0 + b E^0 \langle \varphi_1^0 | \varphi_2^0 \rangle + a H'_{11} + b H'_{12} \\ & = a E^0 + b E^0 \langle \varphi_1^0 | \varphi_2^0 \rangle \end{aligned}$$

$$a(E^0 - E) + a H'_{11} + b H'_{12} = 0 \quad (7)$$

$$\langle \psi_2 | b E^0 + a H'_{21} + b H'_{22} = E b \quad (8)$$

(5)

Solve (7) and (8)

$$\begin{pmatrix} E^0 + H'_{11} - E & H'_{12} \\ H'_{21} & E^0 + H'_{22} - E \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = 0$$

This is an eigenvalue equation which can be solved by diagonalization.

We can write

$$E = E^0 + E'$$

then the eq. for E' is

$$\det \begin{pmatrix} H'_{11} - E' & H'_{12} \\ H'_{21} & H'_{22} - E' \end{pmatrix} = 0$$

Note $H'_{ij} = H'_{ji}$ since H' is Hermitian

We can generalize to n -fold degenerate systems.

The first-order wavefunctions are obtained by solving a and b .

Applications and examples — procedure

- (1) Identify H^0 and H'
- (2) solve eigenstates of H^0 needed
- (3) calculate the matrix elements
→ use symmetry conditions

- (4) finish the calculation.