

Problem 1 Griffiths 1.5

$$\Psi(x, t) = Ae^{-\lambda|x|}e^{-i\omega t}$$

$$\begin{aligned} \text{(a)} \quad 1 &= \int_{-\infty}^{\infty} |\Psi|^2 dx = \int_{-\infty}^{\infty} \Psi \Psi^* dx = A^2 \int_{-\infty}^{\infty} e^{-2\lambda|x|} dx = 2A^2 \int_0^{\infty} e^{-2\lambda x} dx \\ &= 2A^2 \left(\frac{e^{-2\lambda x}}{-2\lambda} \right) \Big|_0^{\infty} = \frac{A^2}{\lambda} \Rightarrow A = \sqrt{\lambda} \end{aligned}$$

$$\begin{aligned} \text{(b)} \quad \langle x \rangle &= \int_{-\infty}^{\infty} x |\Psi|^2 dx = A^2 \int_{-\infty}^{\infty} x e^{-2\lambda|x|} dx = 0 \text{ (odd integrand)} \\ \langle x^2 \rangle &= \int_{-\infty}^{\infty} x^2 |\Psi|^2 dx = A^2 \int_{-\infty}^{\infty} x^2 e^{-2\lambda|x|} dx = 2A^2 \int_0^{\infty} x^2 e^{-2\lambda x} dx \end{aligned}$$

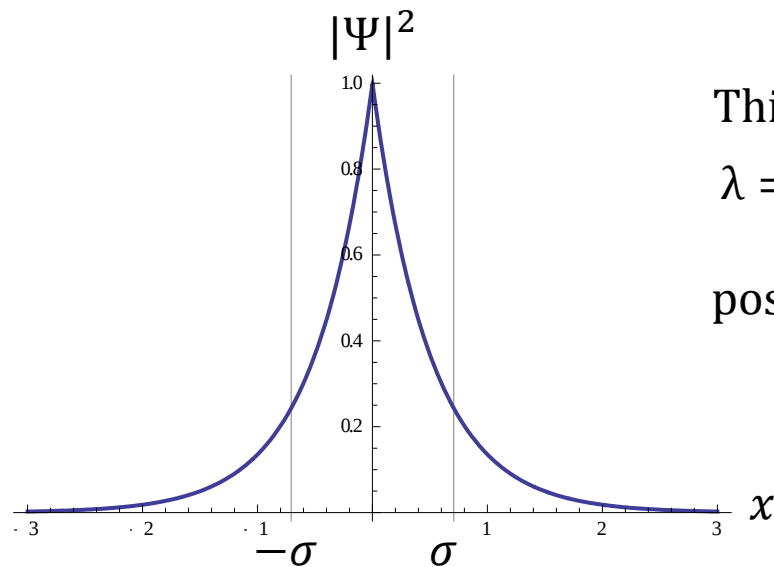
You can do the integration using Mathematica if you don't want to do it by hand. Type the following line in Mathematica then "Shift + Enter":

`Integrate[x^2 * Exp[-2 * λ * x], {x, 0, Infinity}]`

You will get $1/4\lambda^3$

$$\text{So} \quad \langle x^2 \rangle = \frac{2A^2}{4\lambda^3} = \frac{2\lambda}{4\lambda^3} = \frac{1}{2\lambda^2}$$

$$(c) \quad \sigma^2 = \langle x^2 \rangle - \langle x \rangle^2 = \frac{1}{2\lambda^2} \quad \Rightarrow \quad \text{standard deviation } \sigma = \frac{1}{\sqrt{2}\lambda}$$



This figure plots $|\Psi|^2$ as a function of x for $\lambda = 1$. The two vertical lines show the positions of $\sigma \left(= \frac{1}{\sqrt{2}} \right)$ and $-\sigma$.

Probability outside the range $(-\sigma, \sigma)$:

$$\begin{aligned} P_{out} &= \int_{-\infty}^{-\sigma} |\Psi|^2 dx + \int_{\sigma}^{\infty} |\Psi|^2 dx = 2A^2 \int_{\sigma}^{\infty} e^{-2\lambda x} dx = \frac{2A^2}{-2\lambda} e^{-2\lambda x} \Big|_{\sigma}^{\infty} \\ &= e^{-\sqrt{2}} = 0.2431 \end{aligned}$$

Problem 2

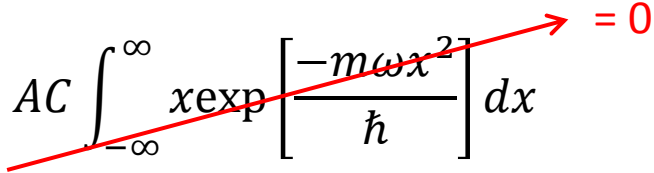
$$(a) \quad H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2$$

$$(b) \quad \varphi_0(x) = A \exp\left[\frac{-m\omega x^2}{2\hbar}\right]$$

$$1 = \int_{-\infty}^{\infty} |\varphi_0|^2 dx = A^2 \int_{-\infty}^{\infty} \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx = A^2 \sqrt{\frac{\pi\hbar}{m\omega}} \Rightarrow A = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4}$$

$$(c) \quad \varphi_1(x) = (B + Cx) \exp\left[\frac{-m\omega x^2}{2\hbar}\right]$$

Orthogonality condition:

$$\begin{aligned} 0 &= \langle \varphi_0 | \varphi_1 \rangle = \int_{-\infty}^{\infty} \varphi_0^* \varphi_1 dx = \int_{-\infty}^{\infty} A(B + Cx) \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx \\ &= AB \int_{-\infty}^{\infty} \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx + AC \int_{-\infty}^{\infty} x \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx \\ &= AB \sqrt{\frac{\pi\hbar}{m\omega}} \Rightarrow B = 0 \end{aligned}$$


Normalization condition:

$$1 = \langle \varphi_1 | \varphi_1 \rangle = \int_{-\infty}^{\infty} |\varphi_1|^2 dx = C^2 \int_{-\infty}^{\infty} x^2 \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx = C^2 \frac{\sqrt{\pi}}{2} \left(\frac{\hbar}{m\omega} \right)^{3/2}$$

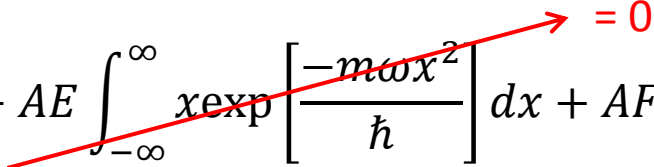
$$\Rightarrow C = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} \sqrt{\frac{2m\omega}{\hbar}}$$

$$\varphi_2(x) = (D + Ex + Fx^2) \exp \left[\frac{-m\omega x^2}{2\hbar} \right]$$

Orthogonality conditions: $\langle \varphi_0 | \varphi_2 \rangle = 0$ & $\langle \varphi_1 | \varphi_2 \rangle = 0$

$$0 = \langle \varphi_0 | \varphi_2 \rangle = \int_{-\infty}^{\infty} \varphi_0^* \varphi_2 dx = \int_{-\infty}^{\infty} A(D + Ex + Fx^2) \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx$$

$$= AD \int_{-\infty}^{\infty} \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx + AE \int_{-\infty}^{\infty} x \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx + AF \int_{-\infty}^{\infty} x^2 \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx$$



$$= AD \sqrt{\frac{\pi\hbar}{m\omega}} + AF \frac{\sqrt{\pi}}{2} \left(\frac{\hbar}{m\omega} \right)^{3/2} \Rightarrow F = -2D \left(\frac{m\omega}{\hbar} \right)$$

$$\begin{aligned}
0 = \langle \varphi_1 | \varphi_2 \rangle &= \int_{-\infty}^{\infty} \varphi_1^* \varphi_2 dx = \int_{-\infty}^{\infty} Cx(D + Ex + Fx^2) \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx \\
&= \cancel{CD \int_{-\infty}^{\infty} x \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx} + CE \int_{-\infty}^{\infty} x^2 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx + \cancel{CF \int_{-\infty}^{\infty} x^3 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx} \\
&= CE \frac{\sqrt{\pi}}{2} \left(\frac{\hbar}{m\omega}\right)^{3/2} \Rightarrow E = 0
\end{aligned}$$

Normalization condition:

$$\begin{aligned}
1 = \langle \varphi_2 | \varphi_2 \rangle &= \int_{-\infty}^{\infty} |\varphi_2|^2 dx = \int_{-\infty}^{\infty} (D + Fx^2)^2 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx \\
&= D^2 \int_{-\infty}^{\infty} \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx + 2DF \int_{-\infty}^{\infty} x^2 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx + F^2 \int_{-\infty}^{\infty} x^4 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx \\
&= D^2 \sqrt{\frac{\pi \hbar}{m\omega}} + 2DF \frac{\sqrt{\pi}}{2} \left(\frac{\hbar}{m\omega}\right)^{3/2} + F^2 \frac{3\sqrt{\pi}}{4} \left(\frac{\hbar}{m\omega}\right)^{5/2} \quad \text{Recall } F = -2D \left(\frac{m\omega}{\hbar}\right) \\
&= 2D^2 \sqrt{\frac{\pi \hbar}{m\omega}} \Rightarrow D = \frac{1}{\sqrt{2}} \left(\frac{m\omega}{\pi \hbar}\right)^{1/4} \quad \text{and} \quad F = -\frac{\sqrt{2}}{\pi^{1/4}} \left(\frac{m\omega}{\hbar}\right)^{5/4}
\end{aligned}$$

$$(d) \quad \varphi_1(x) = Cx \exp \left[\frac{-m\omega x^2}{2\hbar} \right] \quad \text{with } C = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} \sqrt{\frac{2m\omega}{\hbar}}$$

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2 = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2}m\omega^2 x^2$$

$$\begin{aligned} H\varphi_1(x) &= \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2}m\omega^2 x^2 \right) Cx \exp \left[\frac{-m\omega x^2}{2\hbar} \right] \\ &= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \left\{ Cx \exp \left[\frac{-m\omega x^2}{2\hbar} \right] \right\} + \frac{1}{2}m\omega^2 x^2 Cx \exp \left[\frac{-m\omega x^2}{2\hbar} \right] \\ &= -\frac{\hbar^2}{2m} \left\{ -\frac{3m\omega}{\hbar} x + \frac{m^2\omega^2}{\hbar^2} x^3 \right\} C \exp \left[\frac{-m\omega x^2}{2\hbar} \right] \\ &\quad + \frac{1}{2}m\omega^2 x^3 C \exp \left[\frac{-m\omega x^2}{2\hbar} \right] \\ &= \frac{3\hbar\omega}{2} Cx \exp \left[\frac{-m\omega x^2}{2\hbar} \right] \\ &= \frac{3\hbar\omega}{2} \varphi_1(x) \quad \Rightarrow \quad E_1 = \frac{3\hbar\omega}{2} \end{aligned}$$

Problem 3

$$a_+ = \frac{1}{\sqrt{2\hbar m\omega}}(-ip + m\omega x) \qquad a_- = \frac{1}{\sqrt{2\hbar m\omega}}(+ip + m\omega x)$$

$$\begin{aligned}[a_-, a_+] &= \frac{1}{2\hbar m\omega} [ip + m\omega x, -ip + m\omega x] = \frac{1}{2\hbar m\omega} \{im\omega[p, x] - im\omega[x, p]\} \\ &= \frac{1}{2\hbar m\omega} \{im\omega(-i\hbar) - im\omega(i\hbar)\} = 1\end{aligned}$$

$$\begin{aligned}\hbar\omega \left(a_+ a_- + \frac{1}{2} \right) &= \hbar\omega \left[\frac{1}{2\hbar m\omega} (-ip + m\omega x)(ip + m\omega x) + \frac{1}{2} \right] \\ &= \frac{1}{2m} (p^2 + im\omega xp - im\omega px + m^2\omega^2 x^2) + \frac{1}{2} \hbar\omega \\ &= \frac{1}{2m} (p^2 + im\omega[x, p] + m^2\omega^2 x^2) + \frac{1}{2} \hbar\omega \\ &= \frac{1}{2m} (p^2 - \hbar m\omega + m^2\omega^2 x^2) + \frac{1}{2} \hbar\omega \\ &= \frac{p^2}{2m} + \frac{1}{2} m\omega^2 x^2 = H\end{aligned}$$

$$\begin{aligned}
H(a_+\varphi_n) &= \hbar\omega\left(a_+a_- + \frac{1}{2}\right)a_+\varphi_n = \hbar\omega(a_+a_-a_+ + 1/2a_+)\varphi_n \\
&= \hbar\omega a_+\left(a_-a_+ + \frac{1}{2}\right)\varphi_n = \hbar\omega a_+\left(a_+a_- + 1 + \frac{1}{2}\right)\varphi_n \\
&= a_+(H + \hbar\omega)\varphi_n = a_+(E + \hbar\omega)\varphi_n = (E + \hbar\omega)a_+\varphi_n
\end{aligned}$$

Therefore $a_+\varphi_n$ is an eigenstate of H with energy $(E + \hbar\omega)$. But it may not be normalized. It is related to the normalized eigenstate φ_{n+1} up to a coefficient:

$$a_+\varphi_n = C\varphi_{n+1}$$

$$\begin{aligned}
1 &= \langle\varphi_{n+1}|\varphi_{n+1}\rangle = \frac{1}{|C|^2}\langle a_+\varphi_n|a_+\varphi_n\rangle = \frac{1}{|C|^2}\langle\varphi_n|a_-a_+\varphi_n\rangle \\
&= \frac{1}{|C|^2}\langle\varphi_n|(a_+a_- + 1)\varphi_n\rangle = \frac{1}{|C|^2}\langle\varphi_n|(H/\hbar\omega + 1/2)\varphi_n\rangle = \frac{1}{|C|^2}(n+1)
\end{aligned}$$

$$\Rightarrow C = \sqrt{n+1} \quad \text{and} \quad a_+\varphi_n = \sqrt{n+1}\varphi_{n+1} \quad \text{Similarly} \quad a_-\varphi_n = \sqrt{n}\varphi_{n-1}$$

$$\begin{aligned}
H\varphi_n &= \hbar\omega\left(a_+a_- + \frac{1}{2}\right)\varphi_n = \hbar\omega\left(a_+a_-\varphi_n + \frac{1}{2}\varphi_n\right) = \hbar\omega\left(a_+\sqrt{n}\varphi_{n-1} + \frac{1}{2}\varphi_n\right) \\
&= \hbar\omega\left(\sqrt{n}\sqrt{n}\varphi_n + \frac{1}{2}\varphi_n\right) = \hbar\omega\left(n + \frac{1}{2}\right)\varphi_n
\end{aligned}$$

Problem 4

Ground state: $\varphi_0(x) = A \exp\left[\frac{-m\omega x^2}{2\hbar}\right]$ with $A = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4}$

$$\langle x \rangle = \int_{-\infty}^{\infty} \varphi_0^*(x) x \varphi_0(x) dx = A^2 \int_{-\infty}^{\infty} x \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx = 0$$

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} \varphi_0^*(x) x^2 \varphi_0(x) dx = A^2 \int_{-\infty}^{\infty} x^2 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx = A^2 \frac{\sqrt{\pi}}{2} \left(\frac{\hbar}{m\omega}\right)^{3/2} = \frac{\hbar}{2m\omega}$$

$$\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\frac{\hbar}{2m\omega}}$$

$$\begin{aligned} \langle p \rangle &= \int_{-\infty}^{\infty} \varphi_0^*(x) \left(-i\hbar \frac{\partial}{\partial x}\right) \varphi_0(x) dx = A^2 \int_{-\infty}^{\infty} \exp\left[\frac{-m\omega x^2}{2\hbar}\right] (im\omega)x \exp\left[\frac{-m\omega x^2}{2\hbar}\right] dx \\ &= im\omega A^2 \int_{-\infty}^{\infty} x \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx = 0 \end{aligned}$$

$$\langle p^2 \rangle = \int_{-\infty}^{\infty} \varphi_0^*(x) \left(-\hbar^2 \frac{\partial^2}{\partial x^2}\right) \varphi_0(x) dx$$

$$\begin{aligned}
&= -A^2 \hbar^2 \int_{-\infty}^{\infty} \exp\left[\frac{-m\omega x^2}{2\hbar}\right] \left(\frac{m^2\omega^2}{\hbar^2} x^2 - \frac{m\omega}{\hbar}\right) \exp\left[\frac{-m\omega x^2}{2\hbar}\right] dx \\
&= -A^2 m^2\omega^2 \int_{-\infty}^{\infty} x^2 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx + A^2 \hbar m\omega \int_{-\infty}^{\infty} \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx \\
&= -A^2 m^2\omega^2 \frac{\sqrt{\pi}}{2} \left(\frac{\hbar}{m\omega}\right)^{3/2} + A^2 \hbar m\omega \sqrt{\frac{\pi\hbar}{m\omega}} = \frac{1}{2} \hbar m\omega
\end{aligned}$$

$$\sigma_p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \sqrt{\frac{1}{2} \hbar m\omega}$$

$$\sigma_x \sigma_p = \sqrt{\frac{\hbar}{2m\omega}} \sqrt{\frac{1}{2} \hbar m\omega} = \frac{\hbar}{2}$$

Minimum uncertainty

1st excited state:

$$\varphi_1(x) = Cx \exp\left[\frac{-m\omega x^2}{2\hbar}\right] \quad \text{with} \quad C = \frac{\sqrt{2}}{\pi^{1/4}} \left(\frac{m\omega}{\hbar}\right)^{3/4}$$

$$\langle x \rangle = \int_{-\infty}^{\infty} \varphi_1^*(x) x \varphi_1(x) dx = C^2 \int_{-\infty}^{\infty} x^3 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx = 0$$

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} \varphi_1^*(x) x^2 \varphi_1(x) dx = C^2 \int_{-\infty}^{\infty} x^4 \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx = C^2 \frac{3\sqrt{\pi}}{4} \left(\frac{\hbar}{m\omega}\right)^{5/2} = \frac{3\hbar}{2m\omega}$$

$$\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\frac{3\hbar}{2m\omega}}$$

$$\langle p \rangle = \int_{-\infty}^{\infty} \varphi_1^*(x) \left(-i\hbar \frac{\partial}{\partial x}\right) \varphi_1(x) dx$$

$$= -i\hbar C^2 \int_{-\infty}^{\infty} x \left(1 - \frac{m\omega}{\hbar} x^2\right) \exp\left[\frac{-m\omega x^2}{\hbar}\right] dx$$

$$= 0$$

$$\begin{aligned}
\langle p^2 \rangle &= \int_{-\infty}^{\infty} \varphi_1^*(x) \left(-\hbar^2 \frac{\partial^2}{\partial x^2} \right) \varphi_1(x) dx \\
&= -C^2 \hbar^2 \int_{-\infty}^{\infty} \left(\frac{m^2 \omega^2}{\hbar^2} x^4 - \frac{3m\omega}{\hbar} x^2 \right) \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx \\
&= -C^2 m^2 \omega^2 \int_{-\infty}^{\infty} x^4 \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx + 3C^2 \hbar m \omega \int_{-\infty}^{\infty} x^2 \exp \left[\frac{-m\omega x^2}{\hbar} \right] dx \\
&= -\frac{3}{2} \hbar m \omega + 3 \hbar m \omega = \frac{3}{2} \hbar m \omega
\end{aligned}$$

$$\sigma_p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \sqrt{\frac{3}{2} \hbar m \omega}$$

$$\sigma_x \sigma_p = \sqrt{\frac{3\hbar}{2m\omega}} \sqrt{\frac{3}{2} \hbar m \omega} = \frac{3\hbar}{2} > \frac{\hbar}{2}$$

Problem 4: operator approach: MUCH simpler!

$$x = \sqrt{\frac{\hbar}{2m\omega}} (a_+ + a_-) \quad p = i \sqrt{\frac{\hbar m\omega}{2}} (a_+ - a_-)$$

$$\langle n|x|n\rangle = \sqrt{\frac{\hbar}{2m\omega}} \langle n|a_+ + a_-|n\rangle = 0$$

$$\langle n|x^2|n\rangle = \frac{\hbar}{2m\omega} \langle n|a_+a_+ + a_+a_- + a_-a_+ + a_-a_-|n\rangle = \frac{\hbar}{m\omega} (n + 1/2)$$

$$\langle n|p|n\rangle = i \sqrt{\frac{\hbar m\omega}{2}} \langle n|a_+ - a_-|n\rangle = 0$$

$$\langle n|p^2|n\rangle = -\frac{\hbar m\omega}{2} \langle n|a_+a_+ - a_+a_- - a_-a_+ + a_-a_-|n\rangle = \hbar m\omega (n + 1/2)$$

$$\text{Ground state: } \sigma_x \sigma_p = \sqrt{\frac{\hbar}{2m\omega}} \sqrt{\frac{1}{2} \hbar m\omega} = \frac{\hbar}{2}$$

$$\text{1st excited state: } \sigma_x \sigma_p = \sqrt{\frac{3\hbar}{2m\omega}} \sqrt{\frac{3}{2} \hbar m\omega} = \frac{3\hbar}{2}$$

Problem 5

(a) Probability to be in the ground state =

$$|\langle \varphi_1(x) | \psi_A(x, t = 0) \rangle|^2 = \left| \left\langle \varphi_1(x) \left| \frac{1}{\sqrt{2}} \varphi_1(x) + \frac{1}{\sqrt{2}} \varphi_2(x) \right. \right\rangle \right|^2 = \frac{1}{2}$$

(b)
$$\psi_A(x, t) = \frac{1}{\sqrt{2}} \varphi_1(x) e^{-iE_1 t/\hbar} + \frac{1}{\sqrt{2}} \varphi_2(x) e^{-iE_2 t/\hbar}$$

For a 1D infinite square well $E_n = \frac{n^2 \pi^2 \hbar^2}{2ma^2}$, $E_1 = \frac{\pi^2 \hbar^2}{2ma^2}$, $E_2 = \frac{2\pi^2 \hbar^2}{ma^2}$,

$$\psi_A(x, t) = \frac{1}{\sqrt{2}} \varphi_1(x) \exp\left(-\frac{i\pi^2 \hbar t}{2ma^2}\right) + \frac{1}{\sqrt{2}} \varphi_2(x) \exp\left(-\frac{i2\pi^2 \hbar t}{ma^2}\right)$$

(c) Density distribution at $t = 0$ is:

$$|\psi_A(x, t = 0)|^2 = \psi_A^*(x, t = 0) \psi_A(x, t = 0)$$

$$= \left(\frac{1}{\sqrt{2}} \varphi_1^*(x) + \frac{1}{\sqrt{2}} \varphi_2^*(x) \right) \left(\frac{1}{\sqrt{2}} \varphi_1(x) + \frac{1}{\sqrt{2}} \varphi_2(x) \right)$$

$$= \frac{1}{2} (|\varphi_1|^2 + |\varphi_2|^2 + 2\varphi_1\varphi_2) \quad \text{since } \varphi_1 \text{ \& } \varphi_2 \text{ are real}$$

$$\begin{aligned} |\psi_A(x, t)|^2 &= \psi_A^*(x, t)\psi_A(x, t) \\ &= \left(\frac{1}{\sqrt{2}} \varphi_1^*(x) e^{iE_1 t/\hbar} + \frac{1}{\sqrt{2}} \varphi_2^*(x) e^{iE_2 t/\hbar} \right) \left(\frac{1}{\sqrt{2}} \varphi_1(x) e^{-iE_1 t/\hbar} + \frac{1}{\sqrt{2}} \varphi_2(x) e^{-iE_2 t/\hbar} \right) \\ &= \frac{1}{2} (|\varphi_1|^2 + |\varphi_2|^2 + \varphi_1^* \varphi_2 e^{-i(E_2 - E_1)t/\hbar} + \varphi_1 \varphi_2^* e^{i(E_2 - E_1)t/\hbar}) \\ &= \frac{1}{2} (|\varphi_1|^2 + |\varphi_2|^2 + 2\varphi_1 \varphi_2 \cos((E_2 - E_1)t/\hbar)) \end{aligned}$$

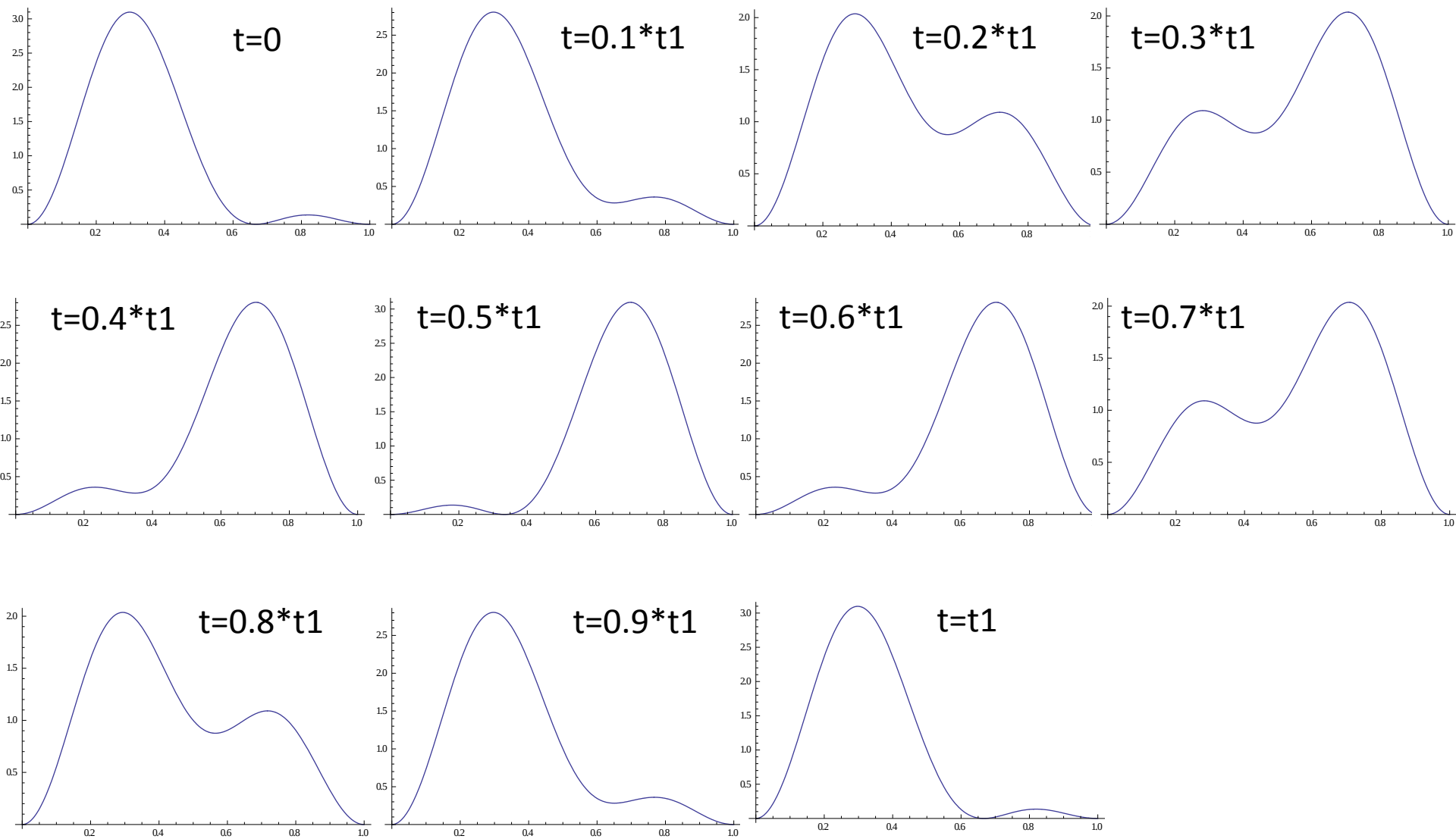
Therefore $|\psi_A(x, t)|^2 = |\psi_A(x, t = 0)|^2$ when $\cos((E_2 - E_1)t/\hbar) = 1$

Or $(E_2 - E_1)t/\hbar = 2\pi, 4\pi, 6\pi, \dots$

$$t = 2n\pi\hbar/(E_2 - E_1) = 2n\pi\hbar / \left[\frac{2\pi^2\hbar^2}{ma^2} - \frac{\pi^2\hbar^2}{2ma^2} \right] = \frac{4nma^2}{3\pi\hbar} \quad \text{where } n = 1, 2, 3, \dots$$

$$\text{The smallest } t_1 = \frac{4ma^2}{3\pi\hbar}$$

(d)



(e)

$$\psi_B(x, t) = \frac{1}{\sqrt{2}} \varphi_1(x) e^{-iE_1 t/\hbar} - \frac{1}{\sqrt{2}} \varphi_2(x) e^{-iE_2 t/\hbar}$$

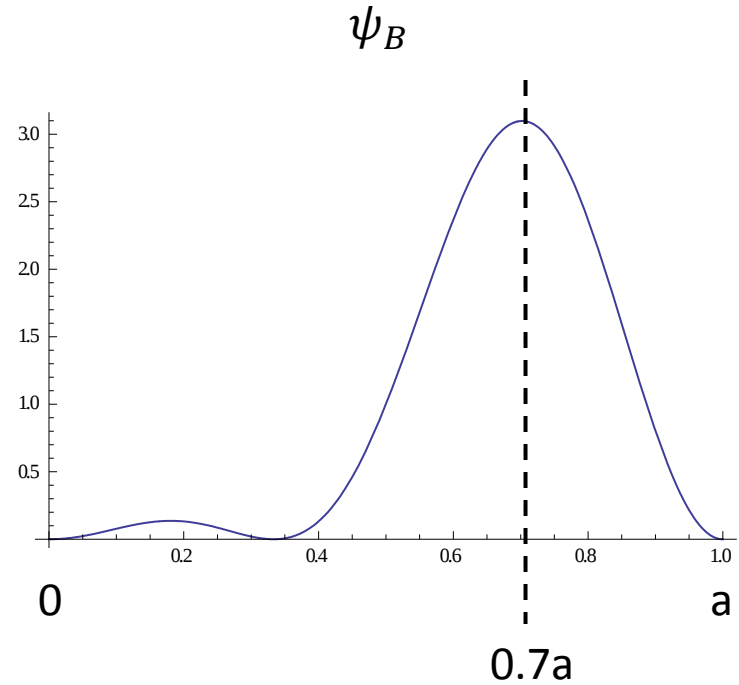
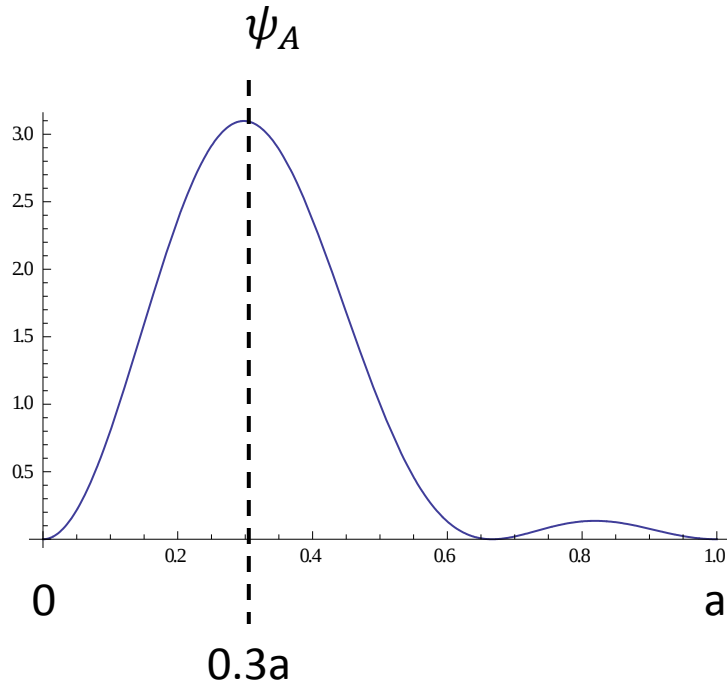
Density distribution at $t = 0$ is:

$$\begin{aligned} |\psi_B(x, t = 0)|^2 &= \psi_B^*(x, t = 0) \psi_B(x, t = 0) \\ &= \left(\frac{1}{\sqrt{2}} \varphi_1^*(x) - \frac{1}{\sqrt{2}} \varphi_2^*(x) \right) \left(\frac{1}{\sqrt{2}} \varphi_1(x) - \frac{1}{\sqrt{2}} \varphi_2(x) \right) \\ &= \frac{1}{2} (|\varphi_1|^2 + |\varphi_2|^2 - 2\varphi_1\varphi_2) \quad \text{since } \varphi_1 \text{ \& } \varphi_2 \text{ are real} \end{aligned}$$

$$\begin{aligned} |\psi_B(x, t)|^2 &= \psi_B^*(x, t) \psi_B(x, t) \\ &= \left(\frac{1}{\sqrt{2}} \varphi_1^*(x) e^{iE_1 t/\hbar} - \frac{1}{\sqrt{2}} \varphi_2^*(x) e^{iE_2 t/\hbar} \right) \left(\frac{1}{\sqrt{2}} \varphi_1(x) e^{-iE_1 t/\hbar} - \frac{1}{\sqrt{2}} \varphi_2(x) e^{-iE_2 t/\hbar} \right) \\ &= \frac{1}{2} (|\varphi_1|^2 + |\varphi_2|^2 - \varphi_1^* \varphi_2 e^{-i(E_2 - E_1)t/\hbar} - \varphi_1 \varphi_2^* e^{i(E_2 - E_1)t/\hbar}) \\ &= \frac{1}{2} (|\varphi_1|^2 + |\varphi_2|^2 - 2\varphi_1 \varphi_2 \cos((E_2 - E_1)t/\hbar)) \end{aligned}$$

t_1 is the same as the previous case

(f)



The two initial states can be determined by measuring the position of the particle:

If the measured value of x is around $0.3a$ then the initial state is ψ_A ;

If the measured value of x is around $0.7a$ then the initial state is ψ_B .