## LETTER TO THE EDITOR

## Reply to "Comment on 'Calculations of energies of intrashell doubly excited states of beryllium-like ions'"

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Abstract. The energy positions of  $1s^23/3l'$  resonances of  $C^{2+}$ ,  $O^{4+}$  and  $Ne^{6+}$  ions are recalculated using two-electron basis functions constructed from eigenstates of the oneelectron model Hamiltonian. The results are in good agreement with the new calculations of Macias *et al* and Bachau *et al*.

In a recent paper, Chen and Lin (1989) calculated the resonance energies of  $1s^23/3l'$  states of  $C^{2+}$ ,  $N^{3+}$ ,  $O^{4+}$  and  $Ne^{6+}$  ions. For some of the lower states, it was pointed out that their results were in disagreement with the previous results of Martin *et al* (1988). In the letter by Bachau *et al* (1990), the discrepancy was attributed to the use of two-electron basis functions built from hydrogenic orbitals adopted by Chen and Lin (1989). It was argued that to project out the open channels appropriately the two-electron basis functions should be constructed from the eigenstates of the one-electron model Hamiltonian.

We have recalculated the energies of these states using two-electron basis functions constructed from the eigenstates of the one-electron Hamiltonian. The results for the two lowest  $1s^23/3l^{-1}S^e$  and  $1s^23/3l'^{-3}P^o$  states are shown in table 1, together with the

**Table 1.** Comparison of the present results with the calculations of Macias *et al* (1989), Chen and Lin (1989) and Martin *et al* (1988) for the energies of  $1s^23/3l'$  states of Be-like C, O and Ne ions. The binding energies are given in eV (1 au = 27.2116 eV).

	States	a	b	c	d	e
C <sup>2+</sup>	(2,0) <sup>1</sup> S <sup>e</sup>	47.37	47.40	47.48	47.74	47.08
	(0,0) <sup>1</sup> S <sup>e</sup>	41.92	41.98	41.86 <sup>f</sup>	41.82	41.91
	(2,0) <sup>3</sup> P <sup>o</sup>	45.98	45.85	45.95	46.12	45.84
	$(0,0)^{3}P^{0}$	40.90	41.18	41.10	40.84	41.18
O <sup>4+</sup>	(2,0) <sup>1</sup> S <sup>e</sup>	107.95			108.33	107.54
	(0,0) <sup>1</sup> S <sup>e</sup>	99.18			99.00	99.18
	$(2,0)^{3}P^{\circ}$	105.73			105.82	105.53
	$(0,0)^{3}P^{\circ}$	97.67			97.58	97.96
Ne <sup>6+</sup>	(2,0) <sup>1</sup> S <sup>e</sup>	193.09	192.74	192.70	193.16	192.28
	(0,0) <sup>1</sup> S <sup>e</sup>	180.78	180.74	180.50	180.46	180.45
	$(2,0)^{3}P^{\circ}$	189.78	189.48	189.62	189.81	189.53
	$(0,0)^{3}P^{\circ}$	178.58	178.98	178.82	178.70	178.99

<sup>a</sup> Present results.

<sup>b</sup> Macias et al (1989), using the pseudopotential-Feshbach approach.

<sup>c</sup> Macias et al (1989), using the conventional-Feshbach approach.

<sup>e</sup> Martin et al (1988).

<sup>f</sup> This state was listed at 41.96 eV in Bachau et al (1990).

<sup>&</sup>lt;sup>d</sup> Chen and Lin (1989).

## L88 Letter to the Editor

newer calculations of Macias *et al* (1989) and Bachau *et al* (1989) and the older ones of Chen and Lin (1989) and Martin *et al* (1988). We first note that the present results are in good agreement with the newer ones of Macias *et al* (1989) and Bachau *et al* (1989). The early larger discrepancy between Chen and Lin (1989) and Martin *et al* (1988) is partly due to the lack of convergence in the results of Martin *et al* (1988). Comparing the difference between the present results and the early ones by us shows that the use of two-electron basis functions constructed from eigenstates of the oneelectron model Hamiltonian only has an effect for lower-Z ions such as  $C^{2+}$ . For the higher Z ions our previous results remain essentially unchanged.

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