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Received 20 Jan 2016 | Accepted 13 May 2016 | Published 22 Jun 2016

DOI: 10.1038/ncomms11922

OPEN

Influence of orbital symmetry on diffraction imaging with rescattering electron wave packets

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The ability to directly follow and time-resolve the rearrangement of the nuclei within molecules is a frontier of science that requires atomic spatial and few-femtosecond temporal resolutions. While laser-induced electron diffraction can meet these requirements, it was recently concluded that molecules with particular orbital symmetries (such as π_g) cannot be imaged using purely backscattering electron wave packets without molecular alignment. Here, we demonstrate, in direct contradiction to these findings, that the orientation and shape of molecular orbitals presents no impediment for retrieving molecular structure with adequate sampling of the momentum transfer space. We overcome previous issues by showcasing retrieval of the structure of randomly oriented O₂ and C₂H₂ molecules, with π_g and π_u symmetries, respectively, and where their ionization probabilities do not maximize along their molecular axes. While this removes a serious bottleneck for laser-induced diffraction imaging, we find unexpectedly strong backscattering contributions from low-Z atoms.

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Knowledge about molecular structure is crucial to the understanding of complex chemical and biological systems. Gas-phase diffraction techniques¹ provide well-established standards to image static structure. However, a deeper understanding of a system and its function only comes with temporally resolved measurements that enable one to observe the nature and timescales of how different components interact with each other. The advent of intense and ultrashort laser pulses prompted the development of techniques such as ultrafast electron diffraction^{2,3} and X-ray diffraction⁴, which are nowadays used to study molecular dynamics on previously inaccessible temporal scales, approaching a few hundred femtoseconds. However, to fully capture the entire dynamics of structural rearrangement, one needs to access the triggering events of a transformation, which demands access to the few-femtosecond temporal range. Not surprisingly, tremendous efforts are being made to access these timescales with current and emerging techniques^{5–8}. Laser-induced electron diffraction (LIED)⁹ is a maturing technique that has already achieved few-femtosecond resolution on homonuclear diatomics,¹⁰ and it has recently been extended to polyatomic molecules¹¹.

In LIED, the momentum distribution of strong-field-induced rescattered electrons is analysed as a function of the incident return momentum (k_r) and scattering angle (θ_r)¹², and spatial resolutions of 5 pm are achievable when combining mid-infrared rescattering¹³ with kinematically complete detection¹¹. While LIED makes use of the doubly differential scattering cross-section as a whole, a simpler analysis was recently presented by Xu *et al.*¹⁴, in which only backscattered ($\theta_r \approx 180^\circ$) electron distributions are analysed. The appeal of using only a one-dimensional (1D) data set instead of the 3D doubly differential scattering cross-section lies in the faster data extraction procedure without the need for iterative algorithms. Moreover, such a Fourier analysis-based 1D version of LIED, FT-LIED, achieves much faster structural retrieval in a manner that is similar to ultrafast electron diffraction. However, when investigating the same simple diatomic species as in ref. 10, Xu *et al.* found that FT-LIED succeeded in retrieving the structure of N₂, but failed for O₂ due to a lack of structural information encoding in the respective backscattered electron distribution. They identified the failure of the FT-LIED methodology as being related to the ionization probability of the target molecule, which maximizes parallel with the bond axis for N₂ but not for O₂. From their findings, and based on their theoretical investigation, the authors concluded that in order to image molecular bonds with the FT-LIED technique the internuclear axis of interest must be aligned with the laser polarization direction. This conclusion has profound implications for LIED as it presents a strong impediment for further advancement of FT-LIED since most larger and complex molecules are not readily aligned (or oriented), and even those that can be would apparently require time-consuming tomographic techniques.

Here, we address the issue of the influence of the orbital symmetry of the target molecule and establish that it does not impede the retrieval of molecular structure. To demonstrate the utility of FT-LIED, we successfully image the structure of both O₂ and C₂H₂ without molecular alignment and irrespective of the fact that molecular ionization is minimal along the molecular axis for both target species. Our results clearly show that the ionization angular dependence of the highest occupied molecular orbital (HOMO) does not prevent imaging of internuclear distances. This demonstrates that imaging is achievable from isotropic gas samples as routinely observed in ultrafast electron diffraction. We do, however, find that molecular structure can be obscured when the momentum transfer space is inadequately sampled. In addition, in our analysis we find unexpectedly

strong scattering from hydrogen atoms. This observation presents a noticeable departure from the theoretical foundation of diffraction imaging, the independent atom model, which underestimates the contribution for the scattering conditions used here and warrants further theoretical investigations for the collision energies accessible to LIED.

Results

LIED imaging of randomly oriented molecules. We first illustrate that the angle of maximum ionization from the HOMO does not need to coincide with the laser polarization axis to retrieve molecular structure. We choose as examples O₂ (blue, solid throughout Fig. 1) and C₂H₂ (red, dashed throughout Fig. 1), whose HOMOs possess π_g and π_u symmetry, respectively, and in neither case do their ionization probabilities maximize parallel to the laser polarization direction (that is, at 0° or 180°), as shown in Fig. 1a. Electrons are predominantly emitted at 90° and 270° for C₂H₂ and at angles $\sim 45^\circ$ away from parallel for O₂. In order not to limit our investigation by the detection method, we resort to 3D electron-ion coincidence detection with a reaction microscope (see Methods for details) to measure the backscattered electron distributions. We have previously shown⁵ that our instrument's coincidence capabilities are important to provide high signal-to-noise data by discriminating backscattered electron distributions from different ionic species that are produced when ionizing the molecular targets. For a simple diatomic this issue is not necessarily a problem as only a few ionic components contribute. In the case of larger and more complex molecules, however, the number of detected ionic fragments will be large and their yields can be comparable to or greater than the main ion¹⁵.

Figure 1b shows the interference signals measured for O₂⁺ and C₂H₂⁺ as a function of the momentum transfer experienced by the diffracted electrons ($q = 2k_r \sin(\theta_r/2)$). We have isolated the molecular modulations from the experimental backscattered electron distributions by implementing empirical background subtraction (details are given in Supplementary Information in the section 'Background subtraction'), in a manner analogous to ultrafast electron diffraction¹. This method presents a number of important advantages compared with previous implementations of LIED. First, the structure of simple molecules can be determined without depending on theoretical calculations. Second, experimental data acquisition times will be reduced significantly as the backscattered electron distributions from partner atoms do not need to be measured. Third, and most importantly, all molecules are now potential targets as even those without partner atoms can be investigated. After background subtraction, it is observed that the resultant interference signals possess a number of extrema and zero-crossings, which are sensitive to molecular structure at the time of rescattering. The phase, frequency and amplitude are observed to be molecule dependent, which is a first indication that structural information is present. We note that, in general, the modulation amplitude of the experimental backscattered electron distributions is greater than that predicted by the independent atom model.

Fourier-transforming the measured interference signals results in a radial distribution spectrum that contains frequency components at the molecular internuclear distances. Before transformation a windowing function and zero padding are applied to the experimental data. The normalized spectra resulting from the Fourier transform of the interference signals in Fig. 1b are presented in Fig. 1c. Also indicated are the expected positions of the O₂ and C₂H₂ cation internuclear distances¹⁶. The inset shows the region between 0.9–1.5 Å, where the anticipated O–O and C–C bond lengths of 1.12 and 1.25 Å, respectively, are located. The experimental peak positions are within 0.05 Å of these values. This finding clearly confirms that

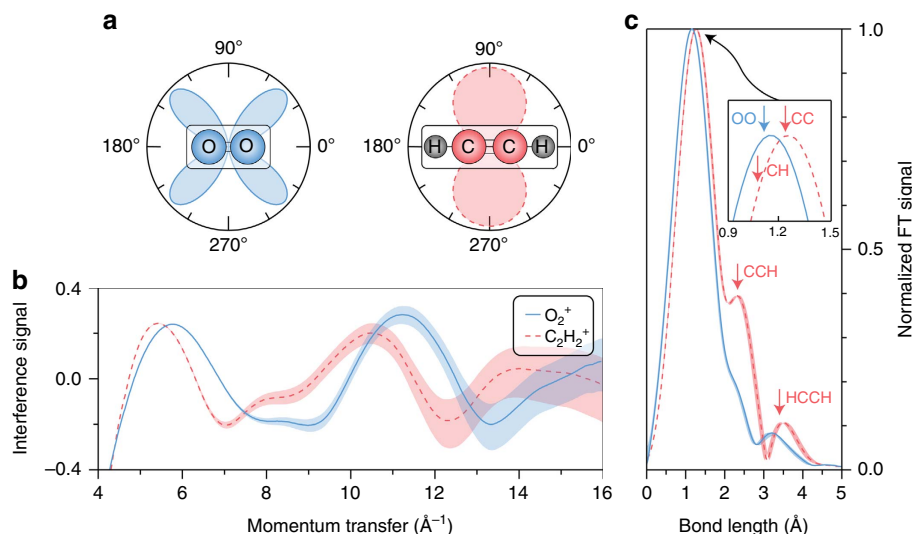


Figure 1 | Extracting bond lengths independent of HOMO structure. (a) Simulated ionization probabilities of the O_2 (blue solid curves throughout figure) and C_2H_2 (red dashed curves throughout figure) HOMOs as a function of the angle between the molecular axes and the laser polarization direction. (b) The interference signals obtained for O_2 and C_2H_2 after ionization of the molecules by the 3.1- μm source. In this panel the shading denotes that the error bars are estimated via Poissonian statistics. (c) The result of Fourier-transforming the interference signals. The expected positions of the O_2 and C_2H_2 cation internuclear distances are indicated. The inset shows a zoomed-in view around 1.2 \AA , where a difference of about 0.14 \AA can be observed between the main peaks of the two molecules. The shaded regions represent the estimated error in the extracted spectra resulting from the uncertainty in the value of the ponderomotive energy.

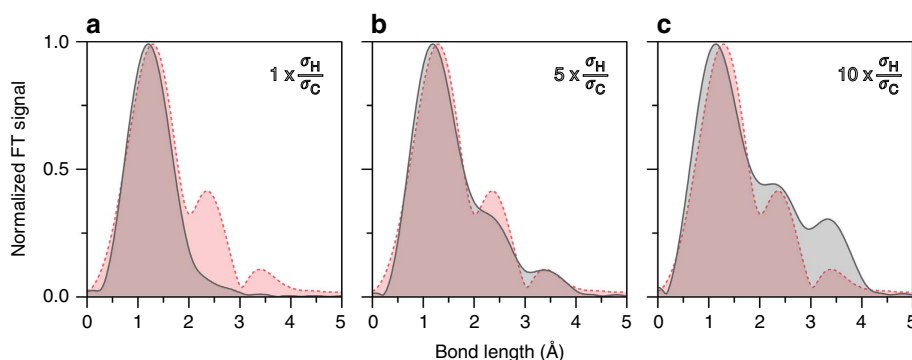


Figure 2 | Limitations of the independent atom model. A comparison of the theoretical (grey, solid throughout figure) and experimental (red, dashed throughout figure) $C_2H_2^+$ Fourier spectra. In (a) the independent atom model is unmodified while in (b) the σ_H/σ_C ratio is increased by a factor of five and in (c) by a factor of ten. All spectra have been normalized to their maximum values.

the bonds of molecules do not need to be aligned with the angle of maximum ionization rate in order to be successfully imaged.

We find that the main C_2H_2 cation peak near 1.3 \AA is likely composed of two internuclear distances that are within <0.2 \AA of each other and therefore not individually resolvable: the C–C bond mentioned above at 1.25 \AA and the C–H (and H–C) bond at 1.08 \AA (see inset of Fig. 1c). Due to the higher cross-section of C compared to H it is expected and observed that this peak should be shifted more towards the position of C–C than C–H. Interestingly and unexpectedly, a second and a third peak are also visible at distances of 2.33 and 3.51 \AA and amplitudes of 0.40 and 0.11, respectively. The position of the second peak agrees with the cation C–C–H and H–C–C internuclear distance of 2.34 \AA . This shows that scattering from the H atom is not negligible as expected from the independent atom model. We discard the third peak of the C_2H_2 spectrum due to low signal to noise even though it contains a signal corresponding to the expected H–C–C–H position. The exclusion of this region is supported by the fact that an artefact is found for O_2 with similarly low amplitude at a similar position (3.21 \AA).

Limitations of the independent atom model. The experimental diffraction data permitted to extract, for both species, O_2 and C_2H_2 , the most prominent contributions of heavy atoms corresponding to O–O and C–C distances, respectively. It is interesting, however, that we obtain clear signatures corresponding to the C–H and C–C–H bond distances. These experimental findings contradict calculations from the independent atom model, which predict that the differential cross-section for H atoms (σ_H) should be one order of magnitude lower than for C atoms (σ_C) for the typical scattering parameters used in LIED ($E_r = 20$ –240 eV and $\theta_r = 180^\circ$). We compare our results to the independent atom model by calculating the expected differential cross-section as a function of returning electron energy at a constant scattering angle of $\theta_r = 180^\circ$. The result of the calculation for neutral C_2H_2 (grey, solid in all panels) is presented alongside our experimental results (red, dashed in all panels) in Fig. 2a. The independent atom model curve is dominated by one frequency component only, which can be observed at a position near 1.21 \AA (neutral C–C bond length) after Fourier transformation. Thus, theory suggests that the influence of H

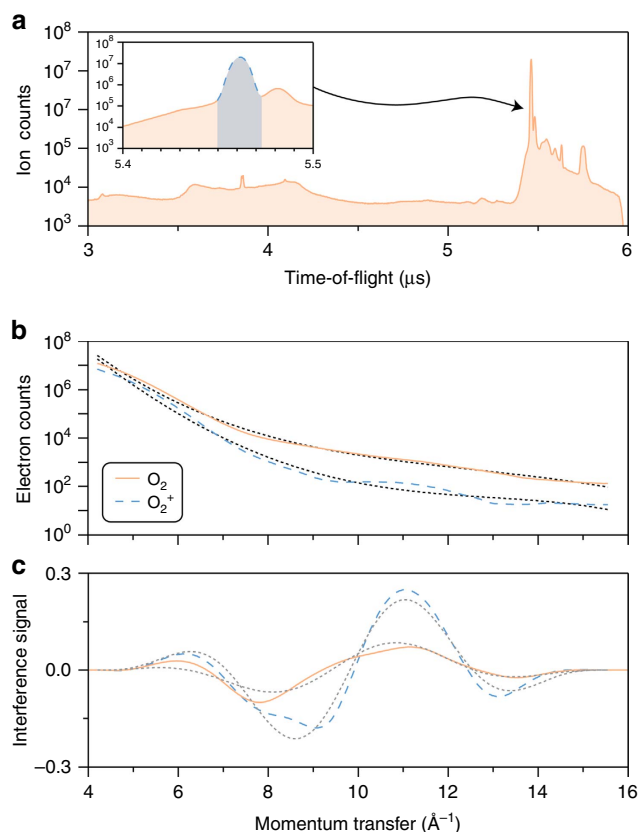


Figure 3 | Importance of correctly sampling the momentum transfer space. (a) The ionic time-of-flight measured after the ionization of O_2 . The inset shows the small temporal range over which the O_2^+ ions are detected (blue shading). Other ions that are not of interest in this manuscript are also detected (orange shading). (b) The extracted backscattered electron distributions when electrons coincident with all ions (orange, solid) or only O_2^+ (blue, dashed) are utilized. Fourth-order polynomial fits are also presented (black, dotted). (c) The corresponding interference signals for O_2^+ (blue, dashed) and all (orange, solid) electrons. Single sinusoidal fits (grey, dotted) to the data show that the observable modulation has decreased by a factor of three.

atoms should not be observable in FT-LIED experiments for our experimentally accessible parameters. Note that the slight difference in the position of the main peak is due to the fact that neutral C_2H_2 is used in the simulation while the cation is measured in the experiment. If the independent atom model calculation is artificially modified by increasing the $\sigma_{\text{H}}/\sigma_{\text{C}}$ ratio by a factor of five, then other frequency components start to contribute to the interference signal. The most interesting finding is that the Fourier-transformed spectrum (Fig. 2b) now closely resembles the experimental one with a new contribution near 2.3 \AA that is similar in magnitude to the measured one. Upon further increasing the $\sigma_{\text{H}}/\sigma_{\text{C}}$ ratio to a factor of ten, three Fourier peaks are visible and the similarity between the two curves clearly shows that the influence of H scattering can be observed in the experimental data (Fig. 2c). The amplitude of the 2.33 \AA peak is about 40% of the main peak and therefore cannot be discarded. Its position is equal to the sum of the retrieved C–C plus C–H bond lengths, therefore corresponding to the distance of C–C–H. Moreover, the smaller Fourier amplitude of C–C–H compared with C–C is consistent with a larger separation, thus indicating that the C–C–H structure was successfully retrieved. We may use identical arguments to state that the small peak near 3.5 \AA is the interference from the two H atoms. Based on the current

independent atom model (Fig. 2a) it is hard to assess whether the experimental peak near 3.5 \AA is related to molecular structure.

Our results suggest that the independent atom model does not fully describe electron scattering for the collision energies that are typical for FT-LIED. We note that differences between C_2H_2 independent atom model calculations and electron impact differential cross-section measurements have also been observed previously for similar scattering energies¹⁷. Using the entire doubly differential cross-section for the data analysis, the independent atom model (IAM) has shown excellent agreement with experimental C_2H_2 LIED data; however, this was at scattering angles $\leq 180^\circ$ (ref. 11). While the apparent failure of the IAM is an important issue that should stimulate further investigations, both theoretical and experimental, our results show that the FT-LIED method continues to work irrespective of this failure. The larger relative contribution of hydrogen with respect to heavier atoms may become an advantage of LIED for studying dynamics involving proton migration, which plays a major role in many chemical reactions.

Importance of sampling the momentum transfer space. We now address why the structure of O_2 could not be imaged in ref. 14. In contrast with ultrafast electron diffraction, LIED makes use of the target molecule's own electrons for imaging. The LIED mechanism is based on strong-field tunnel ionization of the target molecule and rescattering off the target molecular ion. Field-free elastic scattering cross-sections are extracted by invoking the strong-field rescattering model, which can be fulfilled in the quasistatic, or deep tunnelling, regime. The unavoidable consequence of using a strong field to image molecules that possess relatively low ionization potentials (12.07 eV for O_2 (ref. 18)) is the resulting number of strong-field processes such as fragmentation, Coulomb explosion and multiple ionization. Each of these processes has associated rescattering electrons carrying different momenta that obstruct the modulations observed in molecular backscattered electron distributions and therefore prohibit structural retrieval.

To corroborate our explanation we illustrate the influence of these adverse processes on the image bearing backscattered electron distributions. Figure 3a shows ion time-of-flight data after ionization of O_2 from which we had extracted backscattered electron distributions which were only associated to O_2^+ (Fig. 1b). Now we illustrate the importance of our selection of the momentum transfer space by eliminating coincidence conditions. A number of positively charged particles are detected, including the O_2^+ ion located near $5.5 \mu\text{s}$ (blue shading in inset). In Fig. 3b we compare the backscattered electron distributions extracted when using electrons coincident with O_2^+ only (blue, dashed—our conditions) and with all detected ions (orange, solid). Also shown are fourth order polynomial fits to the backscattered electron distributions (black, dotted). Two observations are evident: (1) an order of magnitude difference between the two backscattered electron distributions is visible over a large range of the spectrum, indicating that there are many contaminating electrons even for a simple species such as O_2 , and (2) the superimposed modulation is significantly enhanced when these contaminating electrons are excluded. The factor-of-three difference in modulation is more discernible after subtraction of the empirical background, as presented in Fig. 3c. In this case the signals are displayed with the windowing function already applied. Single sinusoidal fits to the data (grey, dotted) display how well the observed modulations can be represented by a single frequency and also help to quantify the decreased modulation.

The results in Fig. 3 show that even when using a relatively moderate laser intensity of $8.5 \times 10^{13} \text{ W cm}^{-2}$, a significant unwanted electron signal is present during mid-infrared strong-field ionization of molecular targets. In the case of O_2 the desired

electrons cannot be discriminated from this omnipresent unwanted signal without coincidence selection of the momentum transfer space. The peak laser intensity used to investigate O_2 in ref. 14 was over 50% greater than the one used in this study. According to molecular Ammosov-Delone-Krainov simulations¹⁹ this would translate to an order of magnitude higher ionization rate. Thus, even more electrons would be generated from other strong-field processes, and it is these electrons that prohibited structural imaging of O_2 . In the case of a molecule such as N_2 , which was successfully imaged, its higher ionization potential of 15.58 eV (ref. 20) causes it to be harder to both ionize and fragment. Therefore the signal measured would most likely be dominated by N_2^+ electrons.

Discussion

We have investigated whether molecular orientation has a prohibitive influence on the usage of backscattering electron wave packets to image molecular structure. Our investigation provides clear evidence that atomic-scale diffraction imaging does not require alignment of the target species with FT-LIED. Moreover, we use a simple empirically background-subtracted FT-LIED methodology with which we achieve imaging of two entirely different structures, the diatomic O_2 and polyatomic C_2H_2 molecules. Overall, our results have important consequences for FT-LIED and atomic-scale imaging. First, the HOMO structure of an unaligned molecule poses no limitation on diffraction imaging with backscattering wave packets. Second, all molecules, rather than just those that can be aligned to a high degree, will be open to investigation. Lastly, we find that accurate bond lengths, particularly those involving atomic hydrogen, can still be accurately extracted using the FT-LIED method when the IAM begins to fail. This is an important advantage for FT-LIED that makes the investigation of time-resolved studies in complex molecules involving proton migration or isomerization possible. Further theoretical and experimental investigations with other hydrogen-rich molecules will be needed to firmly establish this important aspect of FT-LIED.

Methods

Light source. Our mid-infrared source is a home-built optical parametric chirped pulse amplification system that generates a number of wavelengths including a $\lambda = 3.1 \mu\text{m}$ output that has a duration of 75 fs full width at half maximum at a repetition rate of 160 kHz (ref. 21). The high repetition rate more than compensates for the unfavourable λ^{-4} scaling of strong-field induced rescattering²². A 50-mm-focal-length mirror focuses the $3.1\text{-}\mu\text{m}$ radiation to a spot size of about $6\text{--}7 \mu\text{m}$, which results in a peak intensity of about $I = 8.5 \times 10^{13} \text{ W cm}^{-2}$. The intensity was estimated using a variety of techniques. Such a peak intensity results in a Keldysh parameter of $\gamma \approx 0.3$ and a ponderomotive energy of $U_p = 75 \text{ eV}$, which corresponds to maximum classical return and backscattered electron energies of $E_{\text{ret,max}} = 3.17U_p \approx 240 \text{ eV}$ and $E_{\text{back,max}} = 10U_p \approx 750 \text{ eV}$, respectively. Such high-energy photoelectrons are those that scatter near the nucleus of each atom, where the potential from the core dominates the influence from the outer-shell electron distribution²³.

Detection system. The laser focus intersects with a rotationally cold gas jet ($T \leq 100 \text{ K}$ for both molecules) that passes through two skimmer stages. Both gases are of high purity. The interaction takes place in an ultra-high vacuum chamber that has a base pressure 10^{-11} mbar without gas load and that houses a reaction microscope detection system²⁴. The reaction microscope is able to detect all charged products of the interaction in full coincidence and with 3D momentum information. The magnetic and electric field strengths (12.9 G and 550 V, respectively) were set such that no 'dark regions' (otherwise known as nodes) were located in the half of the electron momentum distribution used during data analysis. There are two reasons why the coincidence capabilities of the reaction microscope are of importance to this study: (1) the capacity to isolate the electrons of interest (i.e., those correlated to O_2^+ or $C_2H_2^+$ only) from those corresponding to other strong-field processes so that uncontaminated backscattered electron distributions can be obtained; and (2) the ability to emulate a time-of-flight spectrometer by extracting backscattered electron distributions that include electronic contributions related to all ion fragments created during the experiment.

Data availability. The data that support the findings of this study are available from the corresponding author upon request.

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Acknowledgements

We acknowledge financial support from the Spanish Ministry of Economy and Competitiveness, through FIS2014-56774-R, FIS2014-51478-ERC, the 'Severo Ochoa' Programme for Centres of Excellence in R&D (SEV-2015-0522), the Catalan Agència de Gestió d'Ajuts Universitaris i de Recerca (AGAUR) with SGR 2014-2016, Fundació Cellex Barcelona, the European Union's Horizon 2020 research and innovation programme under Grant Agreement No. 654148 Laserlab-Europe, the Marie Skłodowska-Curie Grant Agreements No. 641272 and 264666, COST Action MP1203 and COST Action ULIC. M.G.P. is supported by the ICFONEST programme, partially funded by COFUND (FP7-PEOPLE-2013- COFUND) and B.W. was supported by

AGAUR with a PhD fellowship (FI-DGR 2013–2015). A.-T.L. and C.D.L. are supported by the Chemical Sciences, Geosciences, and Biosciences Division, Office of Basic Energy Sciences, Office of Science, U.S. Department of Energy (DOE) under Grant No. DE-FG02-86ER13491.

Author contributions

J.B. conceived the experiment. M.G.P., B.W., M.B., M.S. and H.P. performed the experiment. M.G.P., B.W. and M.S. performed the data analysis. A.-T.L. and C.D.L. provided theoretical support. C.D.S., J.U., R.M. and T.P. provided experimental support. M.G.P. and J.B. wrote the manuscript.

Additional information

Supplementary Information accompanies this paper at <http://www.nature.com/naturecommunications>

Competing financial interests: The authors declare no competing financial interests.

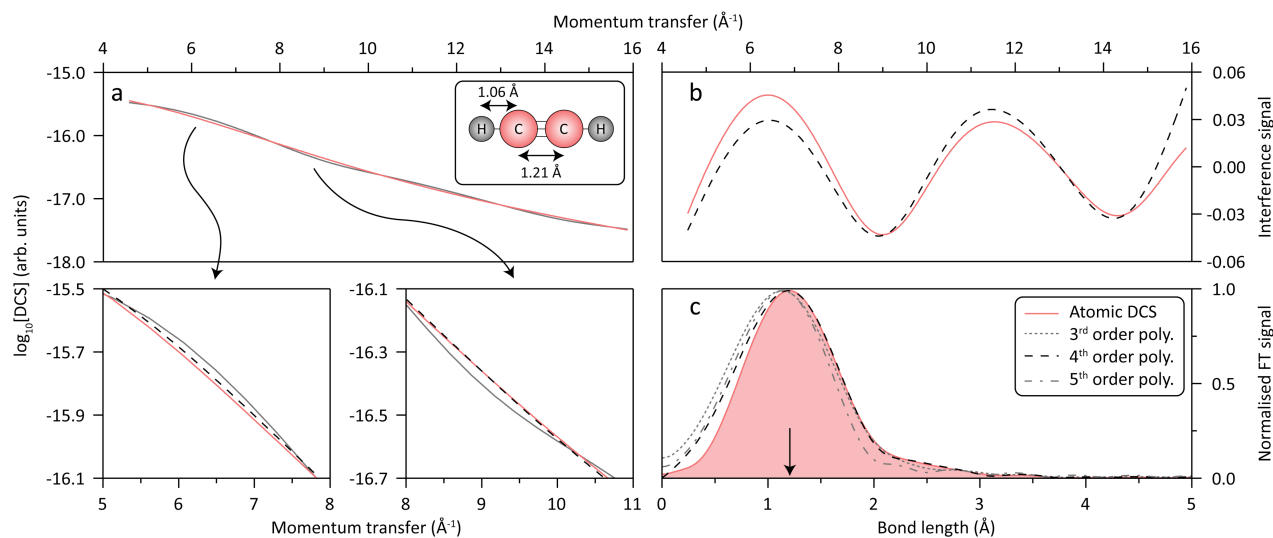
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How to cite this article: Pullen, M. G. *et al.* Influence of orbital symmetry on diffraction imaging with rescattering electron wave packets. *Nat. Commun.* 7:11922 doi: 10.1038/ncomms11922 (2016).

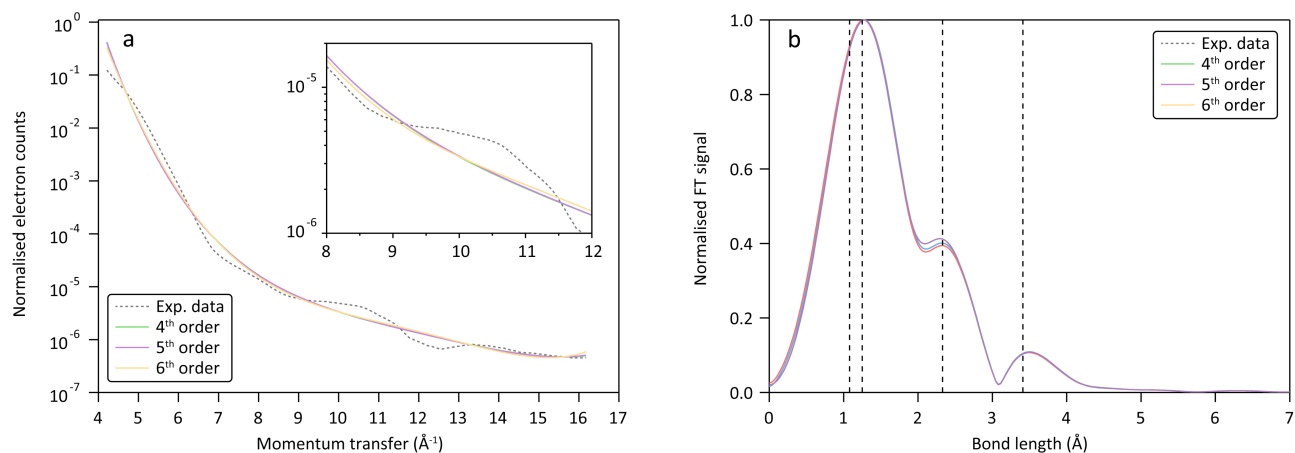


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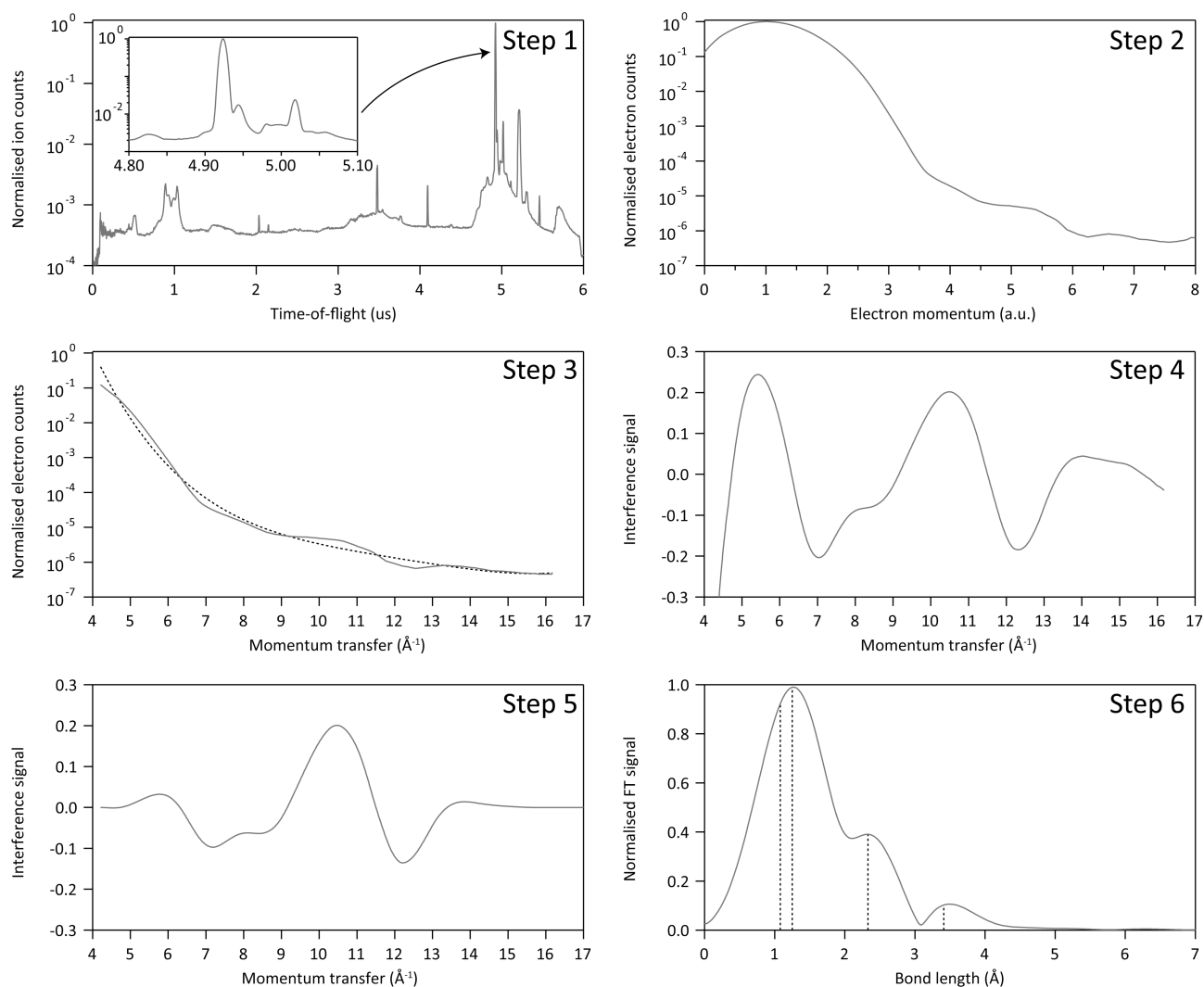
Supplementary Figures



Supplementary Figure 1: Isolation of the interference signal. (a) The calculated molecular (grey solid) and atomic (red solid curves throughout figure) DCSs. The bottom panels show zoomed in regions where modulations of the atomic DCS about the molecular DCS are observed. A fourth order polynomial fit of the molecular DCS (black dashed curves throughout figure) is also shown in each panel. (b) The calculated interference signals (interference signals) after subtraction of the atomic DCS and the polynomial fit. (c) The resultant Fourier transforms of the interference signals alongside the results from third (dotted grey) and fifth (dot-dashed grey) order polynomial fits.



Supplementary Figure 2: Polynomial background subtraction. (a) The detected backscattered electron distribution corresponding to the C_2H_2^+ ion. Polynomial fits of 4th, 5th and 6th order are also presented. The inset shows a zoomed in view between 8 - 12 \AA^{-1} that highlights how the fits lie almost on top of each other. (b) The result of Fourier transforming the interference signals obtained after subtracting the polynomial fits in (a). Very similar behavior is found when investigating O_2 .



Supplementary Figure 3: Data reduction procedure. The steps taken to arrive at the $C_2H_2^+$ bond length spectrum from the experimental data. The same steps are taken to analyse the O2 data. See the text for more details.

Supplementary Discussion

Isolation of molecular information from backscattering electron distributions.

A crucial step in extracting molecular structure imprinted in experimentally measured molecular backscattered electron distributions (ρ_E) the superimposed modulations need to be isolated from the overall distribution by subtracting the monotonically decreasing ‘background’. We therefore outline our procedure before applying it to experimental results. The resulting curve can be called the interference signal and an analogous parameter is used in conventional electron diffraction [1]. In principle, the interference signal can be calculated by subtracting either theoretical simulations or the backscattered electron distributions from atoms with similar ionisation potentials that were measured under the same experimental conditions. Here, and for the first time, we calculate the interference signal from LIED by subtracting an empirically determined background (ρ_B) from the logarithm of ρ_E ,

$$\rho = \log_{10}(\rho_E) - \log_{10}(\rho_B) = \log_{10}(\rho_E/\rho_B).$$

We test our empirical background fitting procedure on independent atom model based simulations that were calculated following procedures that have been outlined previously [11, 23]. Since the independent atom model does not need to take into account the laser-field, we emulate experimental backscattered electron distributions by simply calculating differential cross-sections (DCS) as a function of returning electron energy at a constant scattering angle of $\vartheta_r = 180^\circ$. The calculated molecular DCS (grey, solid) of C_2H_2 is observed to modulate about the corresponding atomic DCS (red, solid) in Supplementary Fig. 1a. These modulations are due to the coherent interference of scattered waves from each atom in the molecule. The bottom two panels show zoomed in regions where the modulations are more visible. The atomic term can be considered equivalent to the background p_B mentioned above. Also shown is a fourth order polynomial fit (black dashed) to the molecular DCS that very closely follows the atomic term. The interference signals resulting from the subtraction of the atomic DCS (red, solid) and the empirical fit (black, dashed) are presented in Supplementary Fig. 1b. The main features of the atomic subtracted interference signal are reproduced by the empirical fit. The normalised Fourier transforms of these signals are presented in Supplementary Fig. 1c alongside the results for when third (grey, dotted) and fifth (grey, dot-dashed) order polynomials are utilised. The most important observation is that the peak of the fundamental frequency component is within 0.05 \AA for all curves. This shows that empirical background subtraction based on polynomial fitting is a robust and accurate method to isolate molecular modulations in FT-LIED. We also tested this polynomial background subtraction procedure on the data presented in Fig. 1a & Supplementary Fig. 3c of Ref. [14] and our results were consistent with those published.

Background subtraction

The robustness of the background subtraction methodology is presented in Supplementary Fig. 2. Thorough tests were performed on the experimental data to ensure that the results of the background fitting algorithm were consistent. In Supplementary Fig. 2a the experimentally measured backscattered electron distribution from C_2H_2 is presented alongside lines of best fit for 4th, 5th and 6th order polynomials. A zoomed in version between $8\text{--}12 \text{ \AA}^{-1}$ shows that these fits lie mostly on top of each other. It was found that these polynomial orders gave almost identical results for both O_2 and C_2H_2 , which highlights the reliability of the background subtraction procedure. The resultant Fourier transformed spectra are presented in Supplementary Fig. 2b where almost identical results are observed. We find that polynomial fits below 3rd order do not adequately remove the ‘background’ signal and fits above 6th order start to remove the real signal.

Experimental data analysis

The procedure to get from the raw experimental data to the bond length spectrum is outlined in Supplementary Fig. 3. We focus on the C_2H_2 case here but the method is the same for O_2 .

Step 1: The ionic time-of-flight shows a variety of positively charged particles detected after single and multiple ionisation as well as fragmentation. The inset shows a zoomed in view of the C_2H_2 cation peak.

Step 2: The electrons associated with the C_2H_2 cation peak are extracted as a function of their final detected momentum. Only electrons that scatter within a $\pm 4^\circ$ cone around the laser polarisation are used in the data analysis.

Step 3: The return momenta of the electrons are calculated from the final momenta using classical trajectory analysis combined with an estimate of the ponderomotive energy of $U_p = 75$ eV. Only the portion of the spectra where modulations are found is used. A polynomial is fitted (dotted curve) to the experimental data (see Supplementary Fig. 3 for more details).

Step 4: The polynomial line of best fit is subtracted from the backscattered electron distribution to isolate the molecular interference signal.

Step 5: A windowing function is applied to the interference signal to reduce spectral leakage and then zero padding is introduced to interpolate in the Fourier domain.

Step 6: After Fourier transformation of the interference signal the bond length spectrum is obtained. Peaks near the expected bond lengths are clearly observed.