Perspectives on the Milestones and Challenges in Gravitational-wave Science: How Did We Get to Today? Beverly K. Berger

LIGO-G1601284



Main focus: Discoveries, inventions, and understanding on the path to gravitational-wave detection from the ground

- Why so long from field equations to detections?
- Why is it hard? EM vs GR
- 4 themes key milestones:
  - Theory (aka mathematical relativity)
  - Simulation (aka numerical relativity)
  - Relativistic astrophysics
  - Experiment: instruments, data analysis
- Convergence: GW150914

General Relativity and Gravitation: a Centennial Perspective, ed. by A. Ashtekar et al, Cambridge U Press (2015), Intro. to Part 2

### The E&M analog:

$$\vec{\nabla} \cdot \vec{D} = \rho$$
  

$$\vec{\nabla} \cdot \vec{B} = 0$$
  

$$\vec{\nabla} \times \vec{H} = \vec{j} + \frac{\partial \vec{D}}{\partial t}$$
  

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$
  
*J. Elide Theorem*

### Field equations published in 1865

James Clerk Maxwell, "A Dynamical Theory of the Electromagnetic Field", *Philosophical Transactions* of the Royal Society of London **155**, 459–512 (1865).

### The E&M analog:



### Hertz experiment performed in 1887



Control EM fields to generate waves and then detect them.

Heinrich Hertz (1893). *Electric Waves: Being Researches on the Propagation of Electric Action with Finite Velocity Through Space*. Dover Publications. ISBN 1-4297-4036-1. English translation

## Field equations presented in 1915

 $R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4}T_{\mu\nu}$ 



688 Sitzung der physikalisch-mathematischen Klasse vom 22. Juni 1916

Näherungsweise Integration der Feldgleichungen der Gravitation.

Von A. Einstein.

Auffectione des Trecoplication die Industrie Pele

Bei der Behandlung der meisten speziellen (nicht prinzipiellen) Probleme auf dem Gebiete der Gravitationstheorie kann man sich damit begnügen, die  $g_{\mu\nu}$  in erster Näherung zu berechnen. Dabei bedient man sich mit Vorteil der imaginären Zeitvariable  $x_4 = it$  aus denselben Gründen wie in der speziellen Relativitätstheorie. Unter »erster Näherung« ist dabei verstanden, daß die durch die Gleichung

 $g_{\mu\nu} = -\delta_{\mu\nu} + \gamma_{\mu\nu}$ 

154 Gesamtsitzung vom 14. Februar 1918. - Mitteilung vom 31. Januar

Über Gravitationswellen.

Von A. EINSTEIN.

#### (Vorgelegt am 31. Januar 1918 [s. oben S. 79].)

Die wichtige Frage, wie die Ausbreitung der Gravitationsfelder erfolgt, ist schon vor anderthalb Jahren in einer Akademiearbeit von mir behandelt worden<sup>1</sup>. Da aber meine damalige Darstellung des Gegenstandes nicht genügend durchsichtig und außerdem durch einen bedauerlichen Rechenfehler verunstaltet ist, muß ich hier nochmals auf die Angelegenheit zurückkommen.

Wie damals beschränke ich mich auch hier auf den Fall, daß das betrachtete zeiträumliche Kontinuum sich von einem »galileischen« nur sehr wenig unterscheidet. Um für alle Indizes

 $g_{\mu\nu} = -\delta_{\mu\nu} + \gamma_{\mu\nu}$ 

(1)

(1)



LIGO detection published in 2016 - detected in 2015

EM	GR
Equations, waves well understood	1st geometric theory, not well understood
Fixed background spacetime	GR provides the spacetime
Detectable laboratory sources	Lab sources <b>not</b> detectable
Strong signals easy to detect	Incredibly weak signals
Room sized experiment	Require kilometer-scale instruments
Cheap	Big science, big bucks
1 author	1000 authors
22 years	100 years

## 1915 – 1918

Linearized GR: gravitational waves

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4}T_{\mu\nu}$$



Einstein, A. 1916. Annalen der Physik, 49, 769–823; 1918. Sitzungsber. Preuss. Akad. Wiss., 154.

$$g_{\mu
u}=\eta_{\mu
u}+h_{\mu
u}$$
etric flat metric pe

perturbation

$$\Box_\eta\,ar{h}_{\mu
u}=0$$
 in vacuum

flat spacetime wave operator

$$\bar{h}_{\mu\nu} = A\epsilon_{\mu\nu}\exp(ik^{\rho}x_{\rho})$$





Linearized GR: gravitational waves **BUT** 

$$\begin{aligned} R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R &= \frac{8\pi G}{c^4}T_{\mu\nu} \\ g_{\mu\nu} &= \eta_{\mu\nu} + h_{\mu\nu} \\ R_{\mu\nu} &= \frac{1}{2}\left(h^{\alpha}_{\mu,\nu\alpha} + h^{\alpha}_{\nu,\mu\alpha} - h_{\mu\nu}, {}^{\alpha}_{\alpha} - h^{\alpha}_{\alpha}, {}^{\mu\nu}\right) \\ \bar{h}_{\mu\nu} &= h_{\mu\nu} - \frac{1}{2}\eta_{\mu\nu}h^{\alpha}_{\alpha} \\ \\ \text{Impose Lorentz gauge: } \bar{h}^{\mu\alpha}, {}_{\alpha} &= 0 \\ \Box_{\eta}\bar{h}_{\mu\nu} &= 0 \\ \\ \text{Are gravitational waves physical?} \end{aligned}$$

Another issue when the linear approximation is not appropriate? How does one separate waves from background?



B.K.B., Ph.D. thesis, 1972

Linearized GR: gravitational wave sources

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4}T_{\mu\nu}$$



Saulson, P. 1994. Fundamentals of interferometric gravitational wave detectors. Singapore: World Scientific

tectors. Singapore: World Scier $\Rightarrow |h| pprox rac{r_S^2}{r_0 r}$ 

reduced moment of inertia at retarded time

$$\bar{h}_{\mu\nu} = \frac{2G}{rc^4} \left( \ddot{\mathcal{I}}_{\mu\nu} \right)_{\rm ret}$$

distance to detector

$$\mathcal{I}_{xx} = 2Mr_0^2 \left(\cos^2 2\pi f_{\rm orb}t - \frac{1}{3}\right)$$

$$h_{xx} = \frac{32\pi^2 G}{rc^4} M r_0^2 f_{\rm orb}^2 \cos 2(2\pi f_{\rm orb})t$$

$$f_{\rm orb} = \frac{GM}{16\pi^2 r_0^3} \qquad r_S = \frac{2GM}{c^2}$$

use Newtonian treatment of orbit

Plug in the numbers: very discouraging in 1918

$$|h| pprox rac{r_S^2}{r_0 r}$$
 For  $1 M_{\odot}$  ,  $r_S pprox 3\,{
m km}$ 

 $r_0 \approx 1 \,\mathrm{AU} \approx 1.5 \times 10^8 \,\mathrm{km}$ Binary stars  $r \approx 100 \,\mathrm{pc} \approx 3 \times 10^{15} \mathrm{km}$  $|h| = \frac{\Delta L}{L} \approx 2 \times 10^{-23}$ ;  $f_{\mathrm{orb}} \approx 1/\mathrm{yr}$ 

Binary neutron stars near merger  $r_0 \approx 20 \,\mathrm{km}$   $r \approx 15 \,\mathrm{Mpc} \approx 4 \times 10^{20} \,\mathrm{km}$  $|h| \approx 10^{-21}$ ;  $f_{\mathrm{orb}} \approx 100 \,\mathrm{Hz}$ 

$$|h| \approx \frac{r_S^2}{r_0 r}$$
 For  $1M_{\odot}$ ,  $r_S \approx 3 \text{ km}$   
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Binary stars  $r \approx 100 \text{ pc} \approx 3 \times 10^{15} \text{ km}$   
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Binary neutron stars near merger  
 $r_0 \approx 20 \text{ km}$ 

 $r \approx 15 \,\mathrm{Mpc} \approx 4 \times 10^{20} \,\mathrm{km}$ 

 $|h| \approx 10^{-21}$ ;  $f_{\rm orb} \approx 100 \,\mathrm{Hz}$ 

Spherically symmetric, vacuum solution in GR:

$$ds^{2} = -\left(1 - \frac{2GM}{c^{2}r}\right)dt^{2} + \frac{dr^{2}}{\left(1 - \frac{2GM}{c^{2}r}\right)} + r^{2}d\Omega^{2}$$
1916

Reduces to Newtonian gravity for field outside mass M for large r and flat spacetime for M = 0,

### BUT

something funny happens at 
$$r=r_S=rac{2GM}{c^2}$$
 .

Physical implications unclear for a long time.

K. Schwarzschild (1916). Sitzungsber. Preuss. Akad. Wiss. 7: 189–196; 18: 424–434.



Karl Schwarzschild

# 1918 - 1955 Not much happened (relevant to this story).

## 1955 - 1965



## Are gravitational waves real? Do they carry energy?

Formalism was developed to identify outgoing radiation (Pirani, Bondi, Sachs, Goldberg, Newman, ...) using null coordinates (light-like trajectories) and the Weyl tensor.



## Are gravitational waves real? Do they carry energy?

Formalism was developed to identify outgoing radiation (Pirani, Bondi, Sachs, Goldberg, Newman, ...) using null coordinates (light-like trajectories) and the Weyl tensor.

Pirani showed that gravitational waves would cause test masses to move.



Sticky bead argument in 1957 demonstrates (theoretically) physical effect of gravitational waves.



via Shane Larson's blog

Bondi, Pirani, Robinson (1959). "Gravitational waves in General Relativity III: Exact plane waves". Proc. Roy. Soc., A 251: 519–533.

Wave makes beads move but friction creates heat. Implies that gravitational waves carry energy.

Schwarzschild solution (1916): behavior at r = 2 M not understood until about 1960, extended to rotating systems in 1963, BH candidates observed in X-rays in 1970s.



Finkelstein, David (1958). "Past-Future Asymmetry of the Gravitational Field of a Point Particle". Phys. Rev 110: 965–967.

1963



## Kerr finds exterior (vacuum) solution for spinning mass



Gravitational Field of a Spinning Mass as an Example of Algebraically Special Metrics Roy P. Kerr, Phys. Rev. Lett. 11, 237 – Published 1 September 1963

# Gravitational wave detection

### J. Weber, schematic of an early bar detector



<1965

Weber, J.; and Wheeler, J.A. (1957). "Reality of the cylindrical gravitational waves of Einstein and Rosen". Rev. Mod. Phys. 29: 509–515. Weber, J., Phys. Rev. Lett. **17**, 1228 (1966).

# Astrophysics

#### 1963



M. Schmidt, Nature 197, 1040 (1963)

Quasars: stellar appearance, cosmological distance



GR taken seriously as relevant to astrophysics. Black holes may exist ????

## 1965 - 1975

# Astrophysics

## Discovery of pulsars





A Hewish, SJ Bell, JDH Pilkington, PF Scott, RA Collins - Nature 217, 709, 1968

# **Uhuru**, an X-ray satellite launched by NASA in 1970, discovered the first evidence for stellar mass black holes.







## Indirect detection of gravitational waves from PSR 1913+16 leads to nobel prize for Hulse and Taylor in 1993.





# Gravitational wave detection

### J. Weber with an early bar detector



Weber, J.; and Wheeler, J.A. (1957). "Reality of the cylindrical gravitational waves of Einstein and Rosen". Rev. Mod. Phys. 29 (3): 509–515.

### J. Weber detection claim in 1969



of Maryland detector coincidence.

J. Weber, Phys. Rev. Lett. 22, 1320 (1969) - Published 16 June 1969
# Computer Simulation

Colliding black hole simulations: U Texas late 1960s



FIG 3. Schematic representation of the coordinate grid for two black holes.

## 1975 - 1990

# Gravitational wave detection



R. Weiss



R. Drever

Interferometric detectors were developed by R. Weiss and R. Drever (in the US). They are scalable and yield waveforms, h(t), rather than energy deposition.



40-meter prototype at Caltech

Unsung hero:



R. Isaacson (Gravitation at NSF)

Rich Isaacson was the Gravity program officer at NSF for 29 years, from 1973.

He convinced his division director and the NSF director to buy into the LIGO project.

He also ensured that early obstacles were overcome.

By the time I arrived at NSF, LIGO was operating smoothly.

R.A. Isaacson, <u>http://ligo.org/magazine/LIGO-magazine-issue-6.pdf#page=14</u>, <u>http://ligo.org/magazine/LIGO-magazine-issue-8.pdf#page=34</u>

Proposal to the National Science Foundation

#### THE CONSTRUCTION, OPERATION, AND SUPPORTING RESEARCH AND DEVELOPMENT OF A

#### LASER INTERFEROMETER GRAVITATIONAL-WAVE OBSERVATORY

Submitted by the CALIFORNIA INSTITUTE OF TECHNOLOGY Copyright © 1989

Rochus E. Vogt Principal Investigator and Project Director California Institute of Technology

Ronald W. P. Drever Co-Investigator California Institute of Technology

Frederick J. Raab Co-Investigator California Institute of Technology Kip S. Thorne Co-Investigator California Institute of Technology

Rainer Weiss Co-Investigator Massachusetts Institute of Technology 1989

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### LIGO realizes this design with test mass mirrors 4 km apart.



GWs from a double neutron star coalescence at a distance of 15 million light years will change LIGO's length by about meters.

The plan in the proposal for LIGO and Advanced LIGO





## 1990 - 2000

# Gravitational wave detection

Progression of sensitivity of Caltech 40-meter convinced me that LIGO would work.



#### **Observational Limit on Gravitational Waves from Binary Neutron Stars in the Galaxy**

B. Allen,<sup>1</sup> J. K. Blackburn,<sup>2</sup> P. R. Brady,<sup>3</sup> J. D. E Creighton,<sup>1,4</sup> T. Creighton,<sup>4</sup> S. Droz,<sup>5</sup> A. D. Gillespie,<sup>2</sup> S. A. Hughes,<sup>4</sup> S. Kawamura,<sup>2</sup> T. T. Lyons,<sup>2</sup> J. E. Mason,<sup>2</sup> B. J. Owen,<sup>4</sup> F. J. Raab,<sup>2</sup> M. W. Regehr,<sup>2</sup> B. S. Sathyaprakash,<sup>6</sup> R. L. Savage, Jr.,<sup>2</sup> S. Whitcomb,<sup>2</sup> and A. G. Wiseman<sup>1</sup>





#### Hanford, WA







## First lock, LIGO Hanford 2 km, 20 Oct 2000



http://www.ligo.caltech.edu/LIGO\_web/firstlock/

# Computer simulations

## The role of simulations (K. Thorne)



1 NOVEMBER 1993

#### **Collision of Two Black Holes**

Peter Anninos,<sup>1</sup> David Hobill,<sup>1,2</sup> Edward Seidel,<sup>1</sup> Larry Smarr,<sup>1</sup> and Wai-Mo Suen<sup>3</sup>

<sup>1</sup>National Center for Supercomputing Applications, Beckman Institute, 405 N. Mathews Avenue, Urbana, Illinois 61801 <sup>2</sup>Department of Physics and Astronomy, University of Calgary, Calgary, Alberta, Canada T2N 1N4 <sup>3</sup>McDonnell Center for the Space Sciences, Department of Physics, Washington University, St. Louis, Missouri 63130 (Received 13 August 1993)

We study the head-on collision of two equal-mass, nonrotating black holes. We consider various cases, from holes surrounded by a common horizon to holes separated by about 20M, where M is the mass of each hole. The wave forms and energy output are computed, showing that normal modes of the final black hole are clearly excited. We also estimate analytically the total gravitational radiation emitted, considering tidal heating of horizons and other effects. The analytic calculations, perturbation theory, and strong-field, nonlinear numerical calculations agree very well with each other.



## Binary black hole simulations (circa 2000):





S. Brandt et al, Phys.Rev.Lett.85:5496-5499,2000

Grazing incidence + excision

Steve Brandt, Randall Correll, Roberto Gómez, Mijan Huq, Pablo Laguna, Luis Lehner, Pedro Marronetti, Richard A. Matzner, David Neilsen, Jorge Pullin, Erik Schnetter, Deirdre Shoemaker, and Jeffrey Winicour

## Why so hard?

Input spacetime grid is not physical; no unique way to write evolution equations - some ill-posed; issues of initial conditions - constraints; .....

Close limit approximation - stopgap until real solution



J. Baker et al, Class. Quantum Grav. 17 No 20 (21 October 2000) L149-L156

# 2000 - today

# Computer simulations

### Computer simulation of orbiting black holes:



2005

**F.** Pretorius

Pretorius used non-standard methods to demonstrate that black hole binaries could be simulated.

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# Binary black hole simulations: standard methods were quickly modified



### UTB/RIT

NASA/GSFC

F. Pretorius, Phys.Rev.Lett.95:121101,2005

M. Campanelli, C. O. Lousto, P. Marronetti, Y. Zlochower, Phys.Rev.Lett. 96:111101,2006

J. G. Baker, J. Centrella, D.-I. Choi, M. Koppitz, J. van Meter, Phys.Rev.Lett. 96:111102,2006

# Gravitational wave detection

# Initial LIGO

#### How sensitive was Initial LIGO?



Approximations for data analysis:

Hundreds of thousands of templates are constructed with approximate waveforms (EOB). These are calibrated with simulations (NR).



A. Buonanno and T. Damour, Phys. Rev. D 59, 084006 (1999)

LIGO Science: The dog that didn't bark in the night: no GW signal from nearby short GRBs rules out compact binary coalescence as origin.



LSC+Hurley, Astrophys. J. 681 (2008) 1419



LSC + others, Astrophys. J. 755 (2012) 2

GRB 051103

GRB 070201

## Advanced LIGO

# The upgrade from LIGO to Advanced LIGO started in 2008 and finished in 2015:



Sampling 1000 X the universe for rare events should increase the event rate by 1000 X.



#### LSC, Class. Quantum Grav. 32 (2015) 074001
#### Advanced LIGO









https://www.advancedligo.mit.edu/

#### January 2013: First lock of Input Mode Cleaner at Hanford



### June 2013: First lock of IR laser using green light in y-arm at Hanford



## October 2013: Aligning installed optics at Livingston

October 2013: Inspecting installed optic at Livingston

# 14 September 2015

#### Advanced LIGO Sensitivity — early October 2015



#### The data:



#### The interpretation:



# A simulation (by SXS group) of GW150914. Many other groups also performed simulations.



#### 図LIGO LIGO Scientific Collaboration





Gravitational waves as solutions of linearized GR:



Schematic of gravitational waves in a non-flat background:

$$g \sim \gamma + h$$
$$\lambda << c \frac{\gamma}{\dot{\gamma}}$$
$$R_{\mu\nu}(\gamma) - \frac{1}{2}\gamma_{\mu\nu}R(\gamma) \sim T_{\mu\nu}^{\text{eff}}[(hh)_{\text{avg}}]$$
$$\Box_{\gamma}h = 0$$

Obtain the background by averaging over the waves.

## Are gravitational waves real? Do they carry energy?

Formalism was developed to identify outgoing radiation (Pirani, Bondi, Sachs, Goldberg, Newman, ...) using null coordinates (light-like trajectories) and the Weyl tensor.

